

# Environmental Chemistry at Vapor/Water Interfaces: Insights from Vibrational Sum Frequency Generation Spectroscopy

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## Keywords

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## Abstract

The chemistry that occurs at surfaces has been an intense area of study for many years owing to its complexity and importance in describing a wide range of physical phenomena. The vapor/water interface is particularly interesting from an environmental chemistry perspective as this surface plays host to a wide range of chemistries that influence atmospheric and geochemical interactions. The application of vibrational sum frequency generation (VSFG), an inherently surface-specific, even-order nonlinear optical spectroscopy, enables the direct interrogation of various vapor/aqueous interfaces to elucidate the behavior and reaction of chemical species within the surface regime. In this review we discuss the application of VSFG to the study of a variety of atmospherically important systems at the vapor/aqueous interface. Chemical systems presented include inorganic ionic solutions prevalent in aqueous marine aerosols, small molecular solutes, and long-chain fatty acids relevant to fat-coated aerosols. The ability of VSFG to probe both the organization and reactions that may occur for these systems is highlighted. A future perspective toward the application of VSFG to the study of environmental interfaces is also provided.

**Interface:** any region between two bulk isotropic media that lacks an inversion center, not just the surface

**Aerosol:** suspended atmospheric particulate matter

**VSFG:** vibrational sum frequency generation

**Beam polarization:** the direction, perpendicular or parallel, of the electric field relative to the incident plane for the associated input or output beam

## 1. INTRODUCTION

The chemistry that occurs at and near environmental interfaces is responsible for a host of important naturally occurring physical phenomena. Examples include transport across biological membranes (1–3), the mobility and fate of aquatic chemical species in the environment (4, 5), and growth and uptake of atmospheric particulate matter (aerosols) (6). As such, the elucidation of environmental interfacial processes has been the goal of many diverse and thorough experimental and theoretical studies; however, a complete understanding of the chemical organization and dynamics that occur for many two-dimensional interfacial regimes is lacking. The advancement of inherently surface-specific nonlinear optical spectroscopies such as second harmonic generation (SHG) spectroscopy and vibrational sum frequency generation (VSFG) spectroscopy allows for the direct interrogation of molecules lacking an inversion center, which naturally occurs at the interface between two bulk isotropic media such as air and water (7–11). The ability to selectively probe interfacial molecules and provide molecular-level information makes techniques such as VSFG powerful tools for the study of interfacial processes. As such, studies utilizing VSFG, along with accompanying theoretical work, have provided much insight into the behavior of molecules at interfaces (10, 12–25).

Although a complete description of VSFG theory is beyond the scope of this review and has been presented elsewhere (21, 26–29), a brief discussion on VSFG theory is necessary for clarity. Conventional VSFG intensity is proportional to the square modulus of the effective second-order nonlinear susceptibility  $|\chi_{eff}^{(2)}|^2$  multiplied by the intensities of the input visible and infrared beams:

$$I_{\text{SFG}} \propto |\chi_{eff}^{(2)}|^2 I_{\text{vis}} I_{\text{IR}} \propto \left| \chi_{eff,\text{NR}}^{(2)} + \sum_{\nu} \chi_{eff,\nu}^{(2)} \right|^2 I_{\text{vis}} I_{\text{IR}}. \quad (1)$$

Here  $I_{\text{SFG}}$ ,  $I_{\text{vis}}$ , and  $I_{\text{IR}}$  are the intensities of the output sum frequency beam, the visible excitation beam, and the infrared excitation beam, respectively, and  $\chi_{eff,\text{NR}}^{(2)}$  and  $\chi_{eff,\nu}^{(2)}$  refer to the effective nonresonant and resonant components of the second-order nonlinear susceptibility.  $\chi_{eff}^{(2)}$  depends on the VSFG experimental setup (i.e., input beam geometry and polarization); most experiments utilize the four most-common beam polarizations: ssp, sps, pss, and ppp, where the first letter corresponds to the output VSFG beam, the second letter corresponds to the input visible beam, and the last letter corresponds to the input infrared beam, with the ssp combination being the most widely utilized. For these polarization combinations, the relationship between  $\chi_{eff}^{(2)}$  and the actual nonlinear susceptibility,  $\chi_{ijk}^{(2)}$ , is given by

$$\chi_{eff,ssp}^{(2)} = L_{yy}(\omega_{\text{VSFG}}) L_{yy}(\omega_{\text{vis}}) L_{zz}(\omega_{\text{IR}}) \sin(\theta_{\text{IR}}) \chi_{yyz}, \quad (2)$$

$$\chi_{eff,sps}^{(2)} = L_{yy}(\omega_{\text{VSFG}}) L_{zz}(\omega_{\text{vis}}) L_{yy}(\omega_{\text{IR}}) \sin(\theta_{\text{vis}}) \chi_{zyy}, \quad (3)$$

$$\chi_{eff,pss}^{(2)} = L_{zz}(\omega_{\text{VSFG}}) L_{yy}(\omega_{\text{vis}}) L_{yy}(\omega_{\text{IR}}) \sin(\theta_{\text{VSFG}}) \chi_{zyy}, \quad (4)$$

$$\begin{aligned} \chi_{eff,ppp}^{(2)} = & -L_{xx}(\omega_{\text{VSFG}}) L_{xx}(\omega_{\text{vis}}) L_{zz}(\omega_{\text{IR}}) \cos(\theta_{\text{VSFG}}) \cos(\theta_{\text{vis}}) \sin(\theta_{\text{IR}}) \chi_{xxx} \\ & -L_{xx}(\omega_{\text{VSFG}}) L_{zz}(\omega_{\text{vis}}) L_{xx}(\omega_{\text{IR}}) \cos(\theta_{\text{VSFG}}) \sin(\theta_{\text{vis}}) \cos(\theta_{\text{IR}}) \chi_{xzx} \\ & +L_{zz}(\omega_{\text{VSFG}}) L_{xx}(\omega_{\text{vis}}) L_{xx}(\omega_{\text{IR}}) \sin(\theta_{\text{VSFG}}) \cos(\theta_{\text{vis}}) \cos(\theta_{\text{IR}}) \chi_{zxx} \\ & +L_{zz}(\omega_{\text{VSFG}}) L_{zz}(\omega_{\text{vis}}) L_{zz}(\omega_{\text{IR}}) \sin(\theta_{\text{VSFG}}) \sin(\theta_{\text{vis}}) \sin(\theta_{\text{IR}}) \chi_{zzz}, \end{aligned} \quad (5)$$

where  $L_{ij}$  is the nonlinear Fresnel factor associated with  $\omega_i$ , and  $\theta_i$  is the input or output angle versus the surface normal for the associated beam (21, 29). By correcting for the nonlinear Fresnel

coefficients, it is possible to remove spectral contributions arising from experimental geometry, as recently demonstrated by Feng et al. (30) for the water spectra at the vapor/liquid interface for spectra originating from many research groups.

The resonant component of the Fresnel factor corrected second-order nonlinear susceptibility is related to the number density,  $N$ , of VSFG active oscillators and the molecular hyperpolarizability,  $\beta_v$ , through the orientationally averaged Euler angle transformation,  $\langle \mu_{IJK:lmn} \rangle$ , between the laboratory coordinates ( $IJK$ ) and the molecular coordinates ( $lmn$ ):

$$\chi_{IJK}^{(2)} = N \sum_{lmn} \langle \mu_{IJK:lmn} \rangle \beta_{lmn}. \quad (6)$$

The molecular hyperpolarizability term (Equation 7) is proportional to the Raman polarizability tensor for the transition moment  $\langle g | \alpha_{lm} | v \rangle$  and the infrared transition moment  $\langle v | \mu_n | g \rangle$ . This gives rise to the VSFG selection rule that a vibrational mode must be both Raman and infrared active for VSFG to be allowed:

$$\beta_{lmn} = \frac{\langle g | \alpha_{lm} | v \rangle \langle v | \mu_n | g \rangle}{\omega_{IR} - \omega_v + i\Gamma_v}, \quad (7)$$

$$\chi_{eff,v}^{(2)} \propto \frac{A_v}{\omega_{IR} - \omega_v + i\Gamma_v}. \quad (8)$$

Generally, however,  $\chi_{eff}^{(2)}$  is what is presented in the majority of published VSFG studies. Equation 8 reveals the Lorentzian lineshape character of the collected VSFG signal, where  $A_v$  is the SFG transition moment strength,  $\omega_v$  is the frequency of the SFG active vibration,  $\omega_{IR}$  is the frequency of the incident infrared laser beam, and  $\Gamma_v$  is the line width of the VSFG transition. It is clear from Equation 1 that during the collection of a conventional VSFG spectrum, the sign of the second-order nonlinear susceptibility is lost. This renders the direct collection of the orientation (phase) of the sum frequency transition impossible. Ji et al. (31) recently developed a generally applicable technique to directly measure the imaginary component (Equation 9) of the second-order nonlinear susceptibility, termed heterodyne-detected phase-sensitive sum frequency generation (PS-SFG), and it is being rapidly adopted by a few researchers within the surface spectroscopy field (32–35):

$$\text{Im}\chi_v^{(2)} = - \sum_v \frac{A_v \Gamma_v}{(\omega_{IR} - \omega_v)^2 + \Gamma_v^2}. \quad (9)$$

The vapor/ neat water interface is perhaps the most simple, common, and important environmental interface. An understanding of this interface is especially vital from an atmospheric chemistry perspective as aqueous atmospheric aerosols are involved in a wide range of phenomena, from heterogeneous chemical reactions to global climate forcing (6, 36, 37). In 1993, Du et al. (11) completed the first study to utilize the VSFG method to investigate neat water structure at the vapor/water interface, showing the existence of three peaks at  $\sim 3,200 \text{ cm}^{-1}$ ,  $\sim 3,400 \text{ cm}^{-1}$ , and  $\sim 3,700 \text{ cm}^{-1}$  within the hydrogen-bonding continuum region between  $3,000$  and  $3,800 \text{ cm}^{-1}$ . The first two broad peaks at  $3,200 \text{ cm}^{-1}$  and  $3,400 \text{ cm}^{-1}$  are attributed to O-H stretching modes of hydrogen-bound water molecules similar to observations of bulk water with Raman and infrared spectroscopies (14), whereas the much narrower peak at  $3,700 \text{ cm}^{-1}$  is assigned to the dangling O-H stretch of water molecules that straddle the vapor/water interface. These dangling water molecules have one O-H bond pointing toward the vapor phase and the other O-H bond pointing toward the bulk liquid, where it is free to hydrogen bond with other water molecules (11). In the years following Du et al.'s study, many other researchers reproduced the VSFG

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**Solute:** a small, soluble, polyatomic molecular species

**Lipid:** molecule featuring a hydrophobic tail such as an alkyl chain with a hydrophilic head group such as a carboxylic acid

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spectrum of the vapor/ neat water interface, yet the exact origin of the VSFG spectral shape for the hydrogen-bonding continuum between 3,000 and 3,600  $\text{cm}^{-1}$  remains controversial (16, 38–43).

Beyond examining the water structure at the vapor/ neat water interface, there is much interest in elucidating molecular behavior via VSFG at the vapor/ water interface for a wide variety of chemical systems. With perspective toward understanding the role that tropospheric aqueous aerosols play in atmospheric chemistry, among other applications, VSFG has been applied to study the behavior that inorganic ions (25, 44–57), small molecular solutes (25, 58–69), and lipids (32, 35, 70–88) exhibit at the vapor/ water interface as well as that of complex solutions that may involve multiple components such as multiple ion species and lipid monolayers spread on ion-containing subphases. At the vapor/ water interface, it is critical to understand these three general classes of chemical systems, which increase in both size and surface preference from simple inorganic ions to solutes to lipids, as they influence the structure, growth, and reactivity of tropospheric aqueous aerosols.

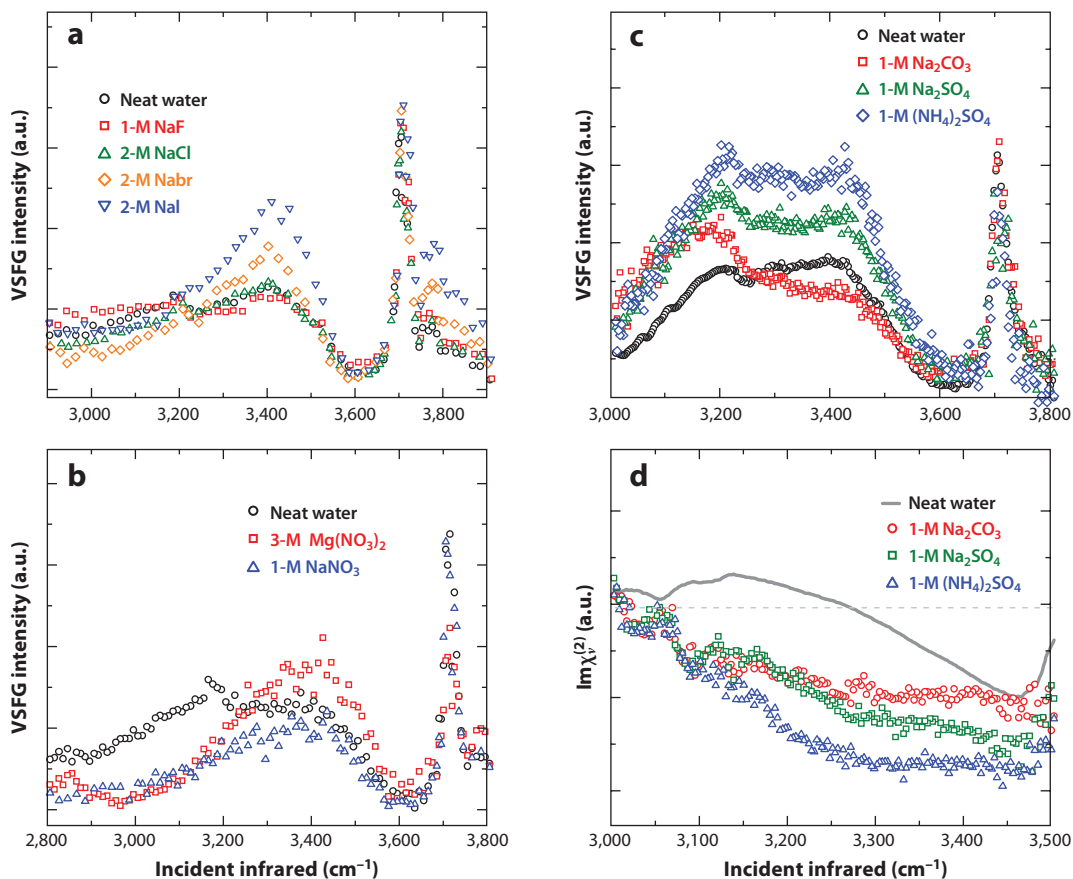
In this review, we present recent insights gained through VSFG studies on the organization and reaction of atmospherically relevant chemical systems (ions, solutes, and lipids) to which the vapor/ water interface plays host. Through the examination of both the response of water's hydrogen-bonding network at the vapor/ liquid interface to the presence of these species and the vibrational modes of the chemical species themselves, elucidation of surface behavior is possible. In Section 2, the surface affinity of various inorganic salt-containing solutions is discussed, along with the effects these ionic species have on interfacial water structure. Included is a brief discussion of the recently developed PS-SFG technique and its application to the study of ions at the vapor/ water interface. Section 3 discusses VSFG results for small molecular aqueous solutes. Here we highlight the surface organization and reaction of methanol at the vapor/ aqueous interface. In Section 4, long-chain fatty acids are discussed, focusing on the surface organization, interaction with ionic species, and oxidation that may occur for these species at vapor/ aqueous interfaces, as these are critical aspects in fat-coated aqueous aerosols. Section 5 briefly discusses future prospects of the application of VSFG to the study of environmental interfaces.

## 2. IONS

### 2.1. Water Structure as Influenced by Ions

The structure of inorganic ions at the vapor/ water interface, as well as the relationship and effect these ions have on the structure of water's hydrogen-bonding network, is of special interest to the atmospheric chemistry community. The use of surface-specific vibrational spectroscopic techniques such as VSFG has been actively pursued by a number of groups to help elucidate long-standing questions on the surface hydrogen-bonding structure of water containing a wide variety of ions (14, 24, 25, 49, 50, 52, 89). Less work has examined the ion modes themselves owing to several factors, ranging from a lack of probable vibrational modes for atomic ions such as halides to the added experimental complexity in producing longer infrared wavelengths necessary to access molecular ionic vibrational modes. However, several recent studies have illustrated that this is possible for a variety of systems, and we discuss these results below (90–92). In this section we provide a brief synopsis of the major findings from work on vapor/ ion-water interfaces using conventional VSFG completed by our lab and others on a number of inorganic salt-containing solutions as well as major unresolved questions remaining for these systems. Recent PS-SFG results on the hydrogen-bonding water structure at the vapor/ aqueous salt solution interface are also highlighted.

Early VSFG work on aqueous inorganic salt solutions at the vapor/ water interface primarily focused on quantifying differences observed between salt solutions and acidic solutions on the



**Figure 1**

Conventional ssp polarized VSGF and  $\text{Im}\chi_v^{(2)}$  SFG spectra of neat water and aqueous solutions of (a) VSGF spectra of sodium halides, (b) VSGF spectra of 3-M  $\text{Mg}(\text{NO}_3)_2$  and 1-M  $\text{NaNO}_3$ , (c) VSGF spectra of 1-M  $\text{Na}_2\text{CO}_3$ , 1-M  $\text{Na}_2\text{SO}_4$ , and 1-M  $(\text{NH}_4)_2\text{SO}_4$ , and (d)  $\text{Im}\chi_v^{(2)}$  SFG spectra of 1-M  $\text{Na}_2\text{CO}_3$ , 1-M  $\text{Na}_2\text{SO}_4$ , and 1-M  $(\text{NH}_4)_2\text{SO}_4$ . Figure adapted from References 46, 57, and 100.

hydrogen-bonding water structure at the vapor/solution interface. In 1997, Raduge et al. (45) and Baldelli et al. (55) were the first to report on the water structure at the vapor/solution interface for sulfuric-acid solutions. Following this work, Baldelli et al. (89) demonstrated in 1999 the influence of alkaline metal salts of sulfate and bisulfate on water structure. Later, Schnitzer et al. (56) extended the discussion to include a wider range of salts versus their acid analog. These experiments and others on a wider range of salts and acids have been repeated by several research groups, including ours (14, 49, 54).

In the past 10 years, the advancement of VSGF instrumentation has renewed interest in the effect that salts, especially halide salts ( $\text{F}^-$ ,  $\text{Cl}^-$ ,  $\text{Br}^-$ , and  $\text{I}^-$ ) with mono- and divalent counter-cations ( $\text{Na}^+$  and  $\text{Mg}^{2+}$ ,  $\text{Ca}^{2+}$ , and  $\text{Sr}^{2+}$ ), have on the water structure near the vapor/liquid interface, prompted by the work of Jungwirth & Tobias (93) predicting surface activity of some halides. The introduction of the halide salts, other than those containing fluoride, caused a slight decrease in the strong hydrogen-bonding region at  $3,200\text{ cm}^{-1}$  and a significant increase in the weak hydrogen-bonding region at  $3,400\text{ cm}^{-1}$  (Figure 1a). These spectral changes increased with increasing polarizability of the anion ( $\text{Cl}^- < \text{Br}^- < \text{I}^-$ ) and were interpreted as an indication

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**Local electric field:**

the electric field generated in the interface by the charge separation of cations and anions in ionic solutions

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of the surface enrichment of solvated anions at the vapor/liquid interface, as an increase in the  $3,400\text{ cm}^{-1}$  region intensity is also a signature for the halide's solvation shell water molecules (46). For fluoride-containing solutions, a slight decrease in both the  $3,200\text{ cm}^{-1}$  and  $3,400\text{ cm}^{-1}$  regions was observed by Raymond & Richmond (54). These observations, along with molecular dynamics simulation results, have been interpreted to represent the  $\text{F}^-$  ion being repelled from the interfacial region (54, 93). In contrast, for halide acid solutions (HCl, HBr, and HI), an enhancement of both the  $3,200\text{ cm}^{-1}$  and  $3,400\text{ cm}^{-1}$  regions is observed with VSFG. For the acids, the increase in the  $3,400\text{ cm}^{-1}$  peak is again attributed to the enrichment of polarizable anions within the interfacial region, whereas the  $3,200\text{ cm}^{-1}$  enhancement is attributed to hydronium ions residing within the interface convoluted with surface potential effects (47). Hydronium ions in halide acid systems have been suggested to both increase the interfacial depth and order water molecules within the interface (47, 50, 94). A comparison of the sodium halide salts to halide acid systems also reveals differences in the free O-H region, with a reduction in the free OH density for halide acids above 1 M (47, 51, 95).

The influence of molecular anions such as  $\text{NO}_3^-$ ,  $\text{SO}_4^{2-}$ , and  $\text{CO}_3^{2-}$  on the hydrogen-bonding structure of water has also been the focus of many VSFG studies (14, 44, 48, 49, 53, 57, 96). An understanding of the VSFG results for polyatomic anion-containing solutions in the water hydrogen-bonding region proves more elusive than that for the halide salt solutions. For all nitrate-containing solutions (**Figure 1b**), a severe depletion of the  $3,200\text{ cm}^{-1}$  peak is observed, and for divalent cation-containing nitrate solutions, an enhancement of the  $3,400\text{ cm}^{-1}$  region is observed, with the larger divalent cation-containing solutions exhibiting a greater enhancement of the  $3,400\text{ cm}^{-1}$  region ( $\text{Sr}^{2+} > \text{Ca}^{2+} > \text{Mg}^{2+}$ ). Solutions containing sulfate, with monovalent cations, however, feature a large enhancement for both the  $3,200\text{ cm}^{-1}$  and  $3,400\text{ cm}^{-1}$  peaks (**Figure 1c**). (Divalent countercation solutions have not been measured.) Finally, sodium-carbonate solutions exhibit an enhancement of the  $3,200\text{ cm}^{-1}$  peak and a depletion in the  $3,400\text{ cm}^{-1}$  peak in VSFG spectra of the vapor/solution interface.

For nitrate systems, the significant depletion of the  $3,200\text{ cm}^{-1}$  peak, and in some cases the enhancement of the  $3,400\text{ cm}^{-1}$  peak, is taken as evidence that the nitrate ions reside within the vapor/water interface. The nitrate ion is thought to have comparable surface activity with that of the chloride ion such that even though nitrate exists within the interfacial region, no enrichment for the ion is observed (52, 57, 93, 96, 97). The magnitude for the decrease and increase of the  $3,200\text{ cm}^{-1}$  and  $3,400\text{ cm}^{-1}$  peaks, respectively, becomes more apparent for divalent cation-containing nitrate solutions as the size of the cation increases ( $\text{Mg}^{2+} < \text{Ca}^{2+} < \text{Sr}^{2+}$ ). This has been attributed to the creation of complex concentration gradients by the larger divalent cations when paired with the nitrate ion such that the depth of the interfacial region increases (57); similar results have been found for the chloride ion when paired with divalent cations (95, 98). Recent PS-SFG studies are consistent with this picture (96).

For sodium-sulfate and ammonium-sulfate solutions (**Figure 1c**), the enhancement for both the  $3,200\text{ cm}^{-1}$  and  $3,400\text{ cm}^{-1}$  regions has been attributed to the ordering of the interfacial water molecules by the local electric field that results from the sulfate dianion residing at the bottom of the interfacial region, with the cations residing some distance above the sulfate dianion (49, 53). Although this interpretation for VSFG results is somewhat speculative, it is in agreement with molecular dynamics simulations and PS-SFG studies for sulfate solutions, as shown in **Figure 1d** (discussed below) (53, 96, 99, 100). The larger spectral enhancement observed for ammonium-sulfate solutions compared with sodium-sulfate solutions is interpreted as a result of the ammonium ions' surface preference, which causes a greater separation between the ammonium and sulfate ions. This increased separation generates a local electric field with a greater magnitude within the interfacial region than what is generated with sodium-sulfate solutions and thus has a greater ordering effect on water molecules within the vapor/solution interface.

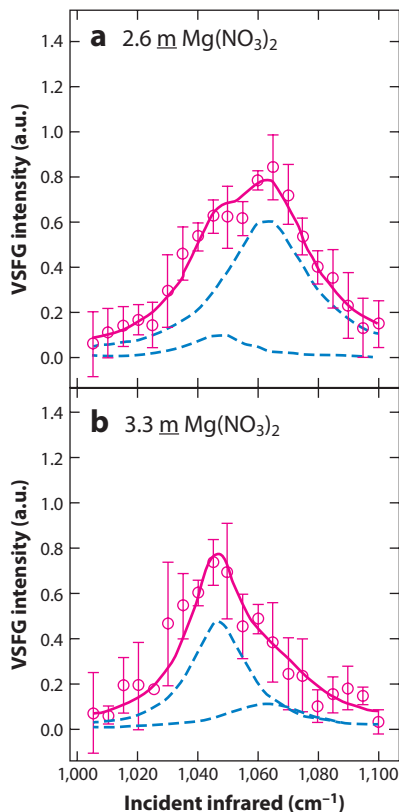
The sodium-carbonate results shown in **Figure 1c** are interpreted as follows: The enhancement of the  $3,200\text{ cm}^{-1}$  peak is generally attributed to a convolution of an increase in the cooperative O-H stretching of tetrahedrally coordinated water molecules, the ordering effects of the ion-induced field, and strong ion-dipole interactions (44, 48, 49). The origin behind the decrease in the  $3,400\text{ cm}^{-1}$  mode for sodium-carbonate solutions is unclear and is the subject of ongoing work.

The recent application of PS-SFG to the study of water organization at vapor/ion-water interfaces has confirmed many interpretations of these interfaces from conventional VSFG studies (48, 96, 100). Shown in **Figure 1d** are the  $\text{Im}\chi_v^{(2)}$  spectra obtained from PS-SFG for sodium-sulfate, ammonium-sulfate, and sodium-carbonate solutions. For both sulfate and carbonate, relative to the neat water surface, the change in sign from positive to negative from  $3,000$  to  $3,200\text{ cm}^{-1}$  and the increased negative intensity of the  $3,200\text{--}3,500\text{ cm}^{-1}$  region for the salt solutions compared to neat water reveal that both sulfate and carbonate have a strong orienting effect on the water OH transition moment, causing the water molecules to align with their OH transition moments pointing toward the bulk solution. This indicates that the dianions are residing near the bottom of the interfacial region, with the associated cations, ammonium and sodium, relatively near the surface. As a first estimate, the differences in spectral magnitude observed for the salt solutions in **Figure 1d** reveal the degree of separation between the cations and anions within the interface for the various salt solutions. Here the sulfate dianion resides deeper within the interface (well below the topmost surface) compared to carbonate, which results in a greater local electric field within the interface for sulfate solutions and exhibits a larger ordering effect on water molecules than is present in carbonate solutions. There is a similar explanation for the differences between ammonium-sulfate and sodium-sulfate solutions, as the ammonium ion is known to exhibit a greater surface preference than sodium, resulting in a larger charge separation of the ion pair within the interface (53).

## 2.2. Ion Interrogation

Beyond the examination of the effect ions have on the water structure at the vapor/water interface, it is possible to utilize VSFG to directly interrogate molecular ions at the vapor/water interface (90–92). These results often reveal surprising differences in ion behavior when compared to bulk studies of aqueous ions utilizing Raman and infrared spectroscopies (92, 101). Although little attention has been given to these types of studies because of added experimental complexity, we feel that several recent studies probing ion modes directly with VSFG at the vapor/water interface should be highlighted. Shown in **Figure 2** are spectra in the N-O vibrational region for the air/aqueous solution interface of  $2.6\text{ m}$  and  $3.3\text{ m}$  magnesium-nitrate solutions (here  $\text{m}$  is moles solute per kilogram solvent). It is also worth noting VSFG spectra published by Miyamae et al. (90) probing the S-O region at the vapor/liquid interface for a series of sulfuric-acid solutions from  $0.01\text{ x}$  to  $0.9\text{ x}$  mole fraction sulfuric acid, along with nitrate-ion spectra from nitric-acid solutions by Soule et al. (91).

The spectra shown in **Figure 2** feature two components that are both attributed to the symmetric stretching mode of the nitrate anion at the vapor/aqueous interface (92). As the concentration increases from  $2.6\text{ m}$  to  $3.3\text{ m}$ , the component at  $\sim 1,047\text{ cm}^{-1}$  increases dramatically. This is taken as evidence of ion-pair formation at the vapor/aqueous interface, in agreement with Raman experiments on the forced dehydration of supersaturated nitrate droplets (102). The work of Miyamae et al. (90) illustrates the surface activity of the sulfate dianion and speciation present as the concentration of  $\text{H}_2\text{SO}_4$  increases from  $0.01\text{ x}$  to  $0.9\text{ x}$  mole fraction  $\text{H}_2\text{SO}_4$ . The little intensity for the dilute  $0.01\text{ x}$   $\text{H}_2\text{SO}_4$  solution indicates that the sulfate dianion is repelled from the interface. However, as the concentration of  $\text{H}_2\text{SO}_4$  increases to  $0.39\text{ x}$ , approximately 98%



**Figure 2**

Conventional ssp polarized VSGF spectra of 2.6  $\underline{m}$   $\text{Mg}(\text{NO}_3)_2$  and 3.3  $\underline{m}$   $\text{Mg}(\text{NO}_3)_2$  solutions (where  $\underline{m}$  is moles solute per kilogram solvent), demonstrating the symmetric stretch of the nitrate ion at the air-aqueous interface. Markers are data, the solid line is the fit, and the dashed lines are fit components. Figure reprinted with permission from Reference 92. Copyright 2009 American Chemical Society.

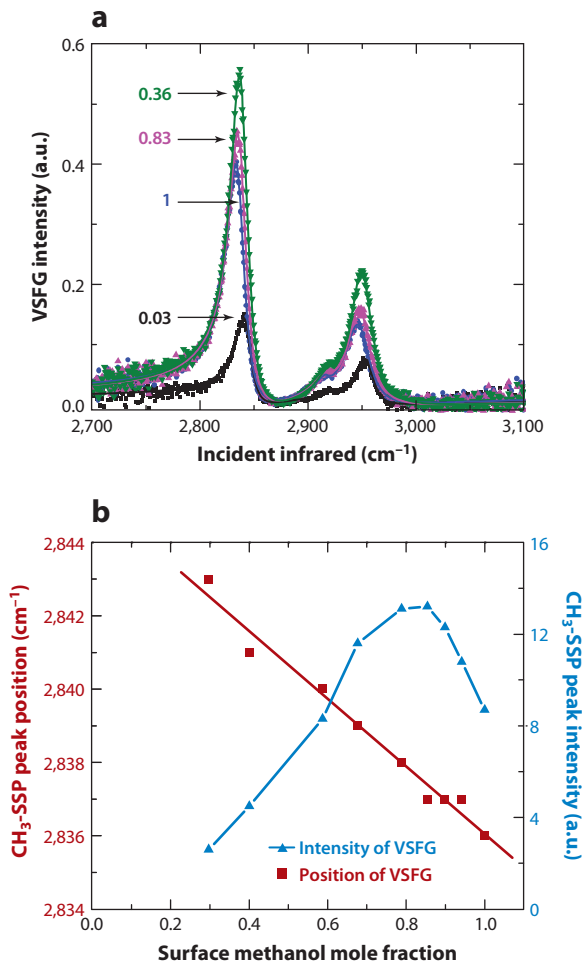
of the sulfate anions are protonated to form the bisulfate,  $\text{HSO}_4^-$ , anion, which can approach the interface. These studies demonstrate the sensitivity of VSGF to atmospherically relevant interfacial processes involving ions (90–92).

### 3. SOLUTES

#### 3.1. Organization of Methanol

The surface preference and reaction of small soluble solutes at the air/vapor interface are also of great interest within the atmospheric chemistry community, and surface-specific optical spectroscopies such as VSGF and PS-SFG have been applied to study a wide range of these chemical systems (25, 58–69). Solutes generally differ in their behavior at vapor/water interfaces from what is observed for ions because of their larger molecular nature and the general, but not absolute, lack of a charge. Here we restrict the discussion to the simple alcohol methanol,  $\text{CH}_3\text{OH}$ , which has broad importance within tropospheric chemistry because it can be used as an alternative fuel and it is a common byproduct of plant biogenesis (6). Methanol is also known to impact both the





**Figure 3**

Conventional ssp polarized VSFG spectra, CH<sub>3</sub>-ssp peak position, and intensity of aqueous methanol (CH<sub>3</sub>OH) solutions. (a) VSFG spectra at different bulk methanol mole fractions. (b) CH<sub>3</sub>-ssp peak position (left y axis) and intensity (right y axis) at different surface methanol mole fractions (with normalized surface number density). Figure adapted from Reference 67.

HO<sub>x</sub> cycle and NO<sub>y</sub> reactions and is involved in various heterogeneous chemical reactions that occur in methanol-containing aqueous aerosols (6, 67, 103–105). We first present the organization of methanol at the vapor/water interface with a perspective toward the evolving interpretation within the literature on its surface behavior. We then highlight the ability of VSFG to monitor surface reactions and the uptake of methanol for various vapor/aqueous interfaces (104). These studies underline the complex nature of reactions that can occur at interfacial regions and provide insight toward heterogeneous chemical reactions that may occur in tropospheric aerosols containing methanol.

Figure 3a shows the conventional VSFG spectra corresponding to a concentration series of binary water-methanol solutions from 0.03 x to 1.0 x mole fraction methanol at the vapor/aqueous interface. The spectra feature two peaks at ~2,840 cm<sup>-1</sup> and ~2,950 cm<sup>-1</sup> and a shoulder on the red side of the 2,950 cm<sup>-1</sup> peak at ~2,920 cm<sup>-1</sup>. The peak at 2,840 cm<sup>-1</sup> is assigned to the

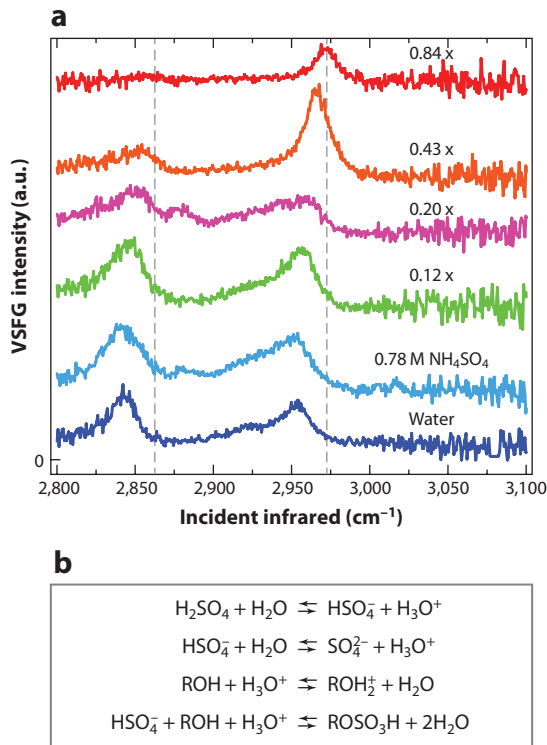
symmetric stretching mode of the methanol  $\text{CH}_3$  moiety, whereas the other peak and shoulder are attributed to the Fermi resonances of the  $\text{CH}_3$  symmetric stretch with the overtones of the  $\text{CH}_3$  bending modes (58, 59, 62, 67, 68). By monitoring the position and intensity of the VSFG  $\text{CH}_3$  symmetric stretch mode, one can infer a picture of methanol's surface behavior. As methanol molecules are surface active compared to water molecules in water-methanol binary mixtures, the bulk mole fraction of methanol does not accurately reflect the surface mole fraction. This has been accounted for in **Figure 3b**, which shows the frequency and intensity shifts versus methanol's surface mole fraction (67). **Figure 3b** illustrates the red shift for the  $\text{CH}_3$  symmetric stretch at the vapor/solution interface observed with VSFG by several researchers (58, 62, 67, 68). This red-shift trend is in agreement with Raman, infrared, and theoretical studies of bulk water-methanol mixtures and is generally accepted as indicating the increased hydrogen-bond-donor character methanol molecules exhibit as the water content decreases (67, 106, 107).

The most striking trend displayed by the spectra in **Figure 3a** is the decrease in the  $\text{CH}_3$  symmetric stretch peak with the increase in methanol concentration above 0.8 x. This trend is not observed for bulk water-methanol studies utilizing Raman and infrared spectroscopies, but it has been observed for all VSFG studies at the vapor/solution interface (58, 62, 67, 68). Early studies by Wolfrum et al. (58) invoked VSFG's intensity dependence on net molecular orientation as well as number density (Equation 6) to explain the decrease in VSFG intensity with increased methanol concentrations. That is, as methanol concentration increases beyond 0.8 x, methanol at the vapor/solution interface becomes more disordered; subsequent VSFG studies by Huang & Wu (68) and Ma & Allen (67) and theoretical calculations by Paul & Chandra (108) confirmed this interpretation.

However, the polarization intensity ratio method (109) used in the above studies to probe methanol's reorientation with increased concentration yielded a large distribution of possible orientations. Recent VSFG studies utilizing the polarization null angle method to determine the molecular orientation for a variety of interfaces have shown the polarization null angle method to be more accurate than the widely used polarization intensity ratio method (21). In applying the polarization null angle methodology to the vapor/methanol-water binary mixture interface, Chen et al. (62) explored the notion that methanol undergoes reorientation with increased concentration at the vapor/solution interface and proposed that this is not the case. Their results indicate that methanol molecules are oriented with the  $\text{CH}_3$  moiety approaching the surface normal directed toward the vapor phase for water-methanol binary mixtures and that this orientation does not significantly change with variation in methanol concentration. To explain the decrease in the VSFG signal observed with an increase in methanol concentration, Chen et al. proposed the presence of an antiparallel structure within the second layer from the surface of the interface, in line with VSFG findings for the vapor/solution interface of acetone-water mixtures (63, 64) and extended X-ray absorption fine-structure studies on the vapor/methanol interface (110). These conclusions have also been supported by the work of Sung et al. (69) for VSFG of alcohol-water binary mixtures and Monte Carlo calculations by Partay et al. (111).

### 3.2. Reactions with Sulfuric Acid

As indicated above for ions, VSFG provides the ability to monitor reactions that solute molecules undergo at surfaces beyond organization. Although many studies of this nature are present in the literature, we focus our discussion on methanol's uptake and reactions with sulfuric-acid solutions at the vapor/solution interface (103–105). Sulfuric acid is the end oxidation product of many sulfur-containing molecules produced in the atmospheric sulfur cycle; sulfur-containing



**Figure 4**

Conventional ssp polarized VSGF spectra and chemical reactions for the uptake of methanol at the surface of 0.12 x to 0.84 x mole fraction H<sub>2</sub>SO<sub>4</sub> solutions. (a) VSGF spectra of methanol surface species at the solution vapor/liquid interfaces after 30 min of exposure to methanol vapor. (b) Relevant chemical reactions for the reaction of methanol with acidic media. Figure reprinted with permission from Reference 104. Copyright 2008 American Chemical Society.

molecules' strong affinity for water enables them to serve as cloud condensation nuclei and thus influence the albedo of the Earth (6, 36, 37).

**Figure 4** illustrates VSGF spectra corresponding to the uptake and reaction of methanol vapor after 30 min of methanol exposure for water and a concentration series of sulfuric-acid solutions, along with the corresponding reaction sequence for the formation of methyl hydrogen sulfate (104). As the concentration of sulfuric acid increases (**Figure 4a**), a blue shift and intensity change are apparent for the two peaks corresponding to the CH<sub>3</sub> symmetric stretch at ~2,840 cm<sup>-1</sup> and the CH<sub>3</sub> Fermi resonance at ~2,954 cm<sup>-1</sup> (compare with the methanol-water mixtures in **Figure 3a**) until the peak at ~2,840 cm<sup>-1</sup> has disappeared, for 0.84 x mole fraction sulfuric-acid solution, and the spectrum shows only one strong peak at ~2,972 cm<sup>-1</sup>. The slight blue shift (~3 cm<sup>-1</sup>) and intensity decrease observed for the 0.12 x and 0.20 x mole fraction sulfuric-acid-methanol solutions are consistent with the formation of protonated methanol, CH<sub>3</sub>OH<sub>2</sub><sup>+</sup>, within the interfacial region via the third reaction in **Figure 4b** (103). In the more concentrated sulfuric-acid solutions, the formation of methyl hydrogen sulfate, CH<sub>3</sub>SO<sub>4</sub>H, occurs within the vapor/solution interface via the fourth reaction pathway in **Figure 4b**; this is apparent from the disappearance of the methanol CH<sub>3</sub> symmetric stretch peak at 2,840 cm<sup>-1</sup> and the appearance of the strong peak at 2,972 cm<sup>-1</sup> attributed to the CH<sub>3</sub> symmetric stretching mode of methyl hydrogen sulfate (103). These findings illustrate the ability of VSGF to monitor volatile organic

compounds such as methanol condensation onto and reaction with aqueous surfaces. Reactions such as the formation of methyl hydrogen sulfate highlight the transformation toward less volatile products, which can contribute to aerosol growth (104).

**PA:** palmitic acid

**OA:** oleic acid

**SAM:** self-assembled monolayer

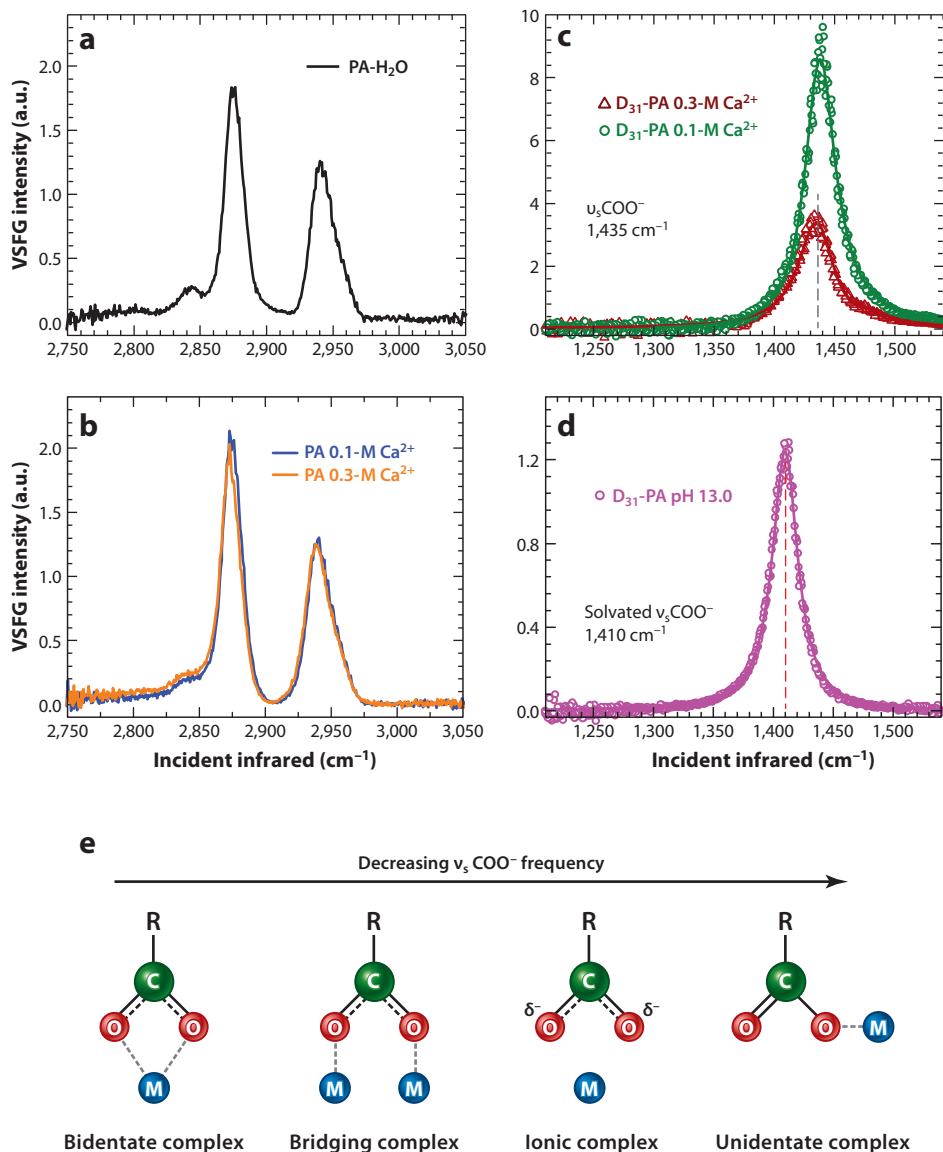
## 4. LIPIDS

### 4.1. Palmitic Acid Organization and Ion Complexation

Organic coated aerosols have been shown to be ubiquitous in both marine and terrestrial environments. Palmitic acid (PA),  $C_{16}H_{32}O_2$ , and oleic acid (OA),  $C_{18}H_{34}O_2$ , are the most prevalent saturated and unsaturated fatty acid organic components found in fat-coated marine aqueous aerosols, respectively (112). Many studies have focused on elucidating the behavior of organic monolayers at the air/aqueous interface with fatty acid Langmuir film proxies, as this behavior has consequences for the growth of fat-coated aerosols and the reactions in which they take part (32, 70, 83, 85, 113–123). VSFG provides an excellent tool for explaining this behavior because of its inherent surface sensitivity. In this section, we review recent VSFG studies on the organization of PA monolayers at the air/aqueous interface as well as examine several reactions that PA monolayers may experience at the air/solution interface, such as ion complexation with ions contained within the solution subphase (74, 86, 88).

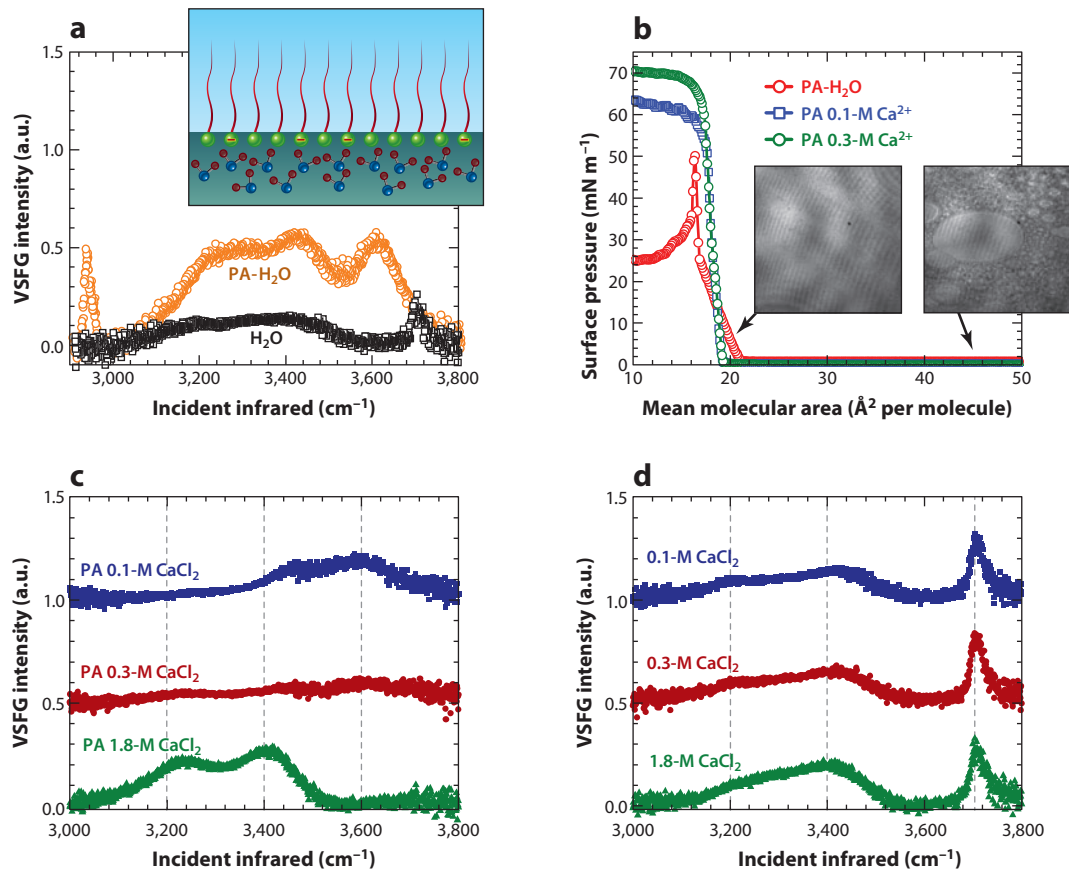
**Figure 5a** shows the spectrum of a PA self-assembled monolayer (SAM) on a neat water subphase in the C-H stretching region for a surface pressure of  $10 \text{ mN m}^{-1}$ . This spectrum features four signatures at  $2,842 \text{ cm}^{-1}$ ,  $2,874 \text{ cm}^{-1}$ ,  $2,940 \text{ cm}^{-1}$ , and  $2,960 \text{ cm}^{-1}$  that have been attributed to the methylene symmetric stretch, the methyl symmetric stretch, the methyl Fermi resonance, and the methyl asymmetric stretch, respectively (74, 88). The relative strength of the methyl symmetric stretch peak at  $2,874 \text{ cm}^{-1}$  compared with the methylene symmetric stretch at  $2,842 \text{ cm}^{-1}$  indicates the high degree of conformational order found in the PA SAM at a surface pressure of  $10 \text{ mN m}^{-1}$ . This is apparent when we consider that for a vibrational mode to be VSFG active, it must not feature an inversion center. The high methyl-to-methylene symmetric stretch ratio indicates that the PA molecules are highly aligned, inducing centrosymmetry between adjacent  $\text{CH}_2$  moieties (124, 125); this feature is observed only for all-*trans* alkyl SAMs with an even number of methylene groups (88, 126).

When a PA monolayer is spread on aqueous subphases that contain ions (0.1- and 0.3-M  $\text{CaCl}_2$ ; **Figure 5b**), VSFG reveals slightly different behavior for the lipid SAM than what is observed for a neat water subphase. Here the overall spectral intensity in the C-H region observed for a PA monolayer is slightly greater than that observed for a PA monolayer on neat water; this suggests either that the PA molecules are more highly aligned when spread on an ionic subphase or that the monolayer packing has increased such that there is a greater number density of PA molecules within the VSFG probe spot. This can also be observed in the compression isotherm (**Figure 6b**) for PA monolayers spread on neat water and  $\text{CaCl}_2$ -containing subphases. To quantitate these observations in the C-H region, Tang and colleagues (74, 88) investigated the COOH head group of PA to discover the effect ions have on the carboxyl and carboxylate symmetric stretching modes present at the vapor/solution interface. Shown in **Figure 5c,d** are spectra corresponding to the PA monolayer head group for PA SAMs spread on a pH-13 aqueous subphase, a 0.1-M  $\text{CaCl}_2$  solution subphase, and a 0.3-M  $\text{CaCl}_2$  subphase (**Figure 5c**). The peak at  $\sim 1,410 \text{ cm}^{-1}$  for the PA monolayer on the pH-13 subphase in **Figure 5d** has been assigned to the solvated  $\text{COO}^-$  symmetric stretch, indicating that the PA monolayer's head group is completely deprotonated at this pH, as is expected from attenuated total reflectance Fourier transform infrared studies by Gershevitc & Sukenik (113) on the pKa for carboxylate-terminated SAM. The spectra in **Figure 5c**



**Figure 5**

Conventional ssp polarized VSGF spectra of palmitic acid (PA) monolayers and schematic representation of ion complexation of the PA head group. (a) VSGF spectra in the C-H stretching region of PA monolayers on neat water at 10 mN m<sup>-1</sup> and near neutral pH. (b) VSGF spectra in the C-H stretching region of a PA monolayer on 0.1- and 0.3-M CaCl<sub>2</sub> solutions. (c) VSGF spectra in the COO<sup>-</sup> stretch region of D<sub>31</sub>-PA monolayers on 0.1- and 0.3-M CaCl<sub>2</sub> solutions. (d) VSGF spectra in the COO<sup>-</sup> stretch region of D<sub>31</sub>-PA monolayers on water with pH 13.0. (e) Illustrations of four possible metal-carboxylate complexes in order of decreasing  $\nu_s \text{COO}^-$  frequency. R represents an alkyl chain. Figure adapted from Reference 74.



**Figure 6**

Conventional ssp polarized VSFG spectra and surface pressure–area isotherms ( $\pi$ - $A$ ) of a palmitic acid (PA) monolayer on aqueous surfaces at pH 6. (a) ssp VSFG spectrum of neat water and the PA monolayer on water at 23°C in the O-H stretching region. (Inset) A schematic of the PA monolayer organization on water. (b) Surface pressure–area isotherms ( $\pi$ - $A$ ) of the PA monolayer on an aqueous surface. (Inset) Brewster angle microscopy images corresponding to PA monolayers at the indicated surface pressures (arrows). (c) VSFG spectra in the O-H stretching region of neat water and CaCl<sub>2</sub> solutions (0.1, 0.3, and 1.8 M). (d) VSFG spectra in the O-H stretching region of PA monolayers on neat water and aqueous CaCl<sub>2</sub> solution (0.1, 0.3, and 1.8 M) subphases. Figure adapted from Reference 86.

corresponding to the PA SAM on CaCl<sub>2</sub> solutions show very different results; here the spectra are dominated by a strong peak at 1,435 cm<sup>-1</sup>, and a slight shoulder at 1,475 cm<sup>-1</sup> was also observed. These spectral features indicate the ion complexation that occurs between the aqueous cation and the ion-induced deprotonated COO<sup>-</sup> head group with greater degrees of complexation, as illustrated in **Figure 5e**, resulting in a larger blue shift in the COO<sup>-</sup> frequency (127–129). The water structure near a PA SAM at the vapor/water and vapor/aqueous salt solution interfaces has also been examined with VSFG to elucidate the ultimate influence exerted by ion complexation and concomitant PA head-group deprotonation (86).

With regard to the water structure at the vapor/water interface for a variety of lipid SAMs, investigators have widely reported that the head-group charge and the packing ability of the SAM are the most dominant forces in influencing water organization near organic monolayers (32). **Figure 6a** illustrates VSFG spectra in the O-H stretching region corresponding to neat water and a PA monolayer on a neat water subphase. The spectral enhancement of the O-H stretching peaks

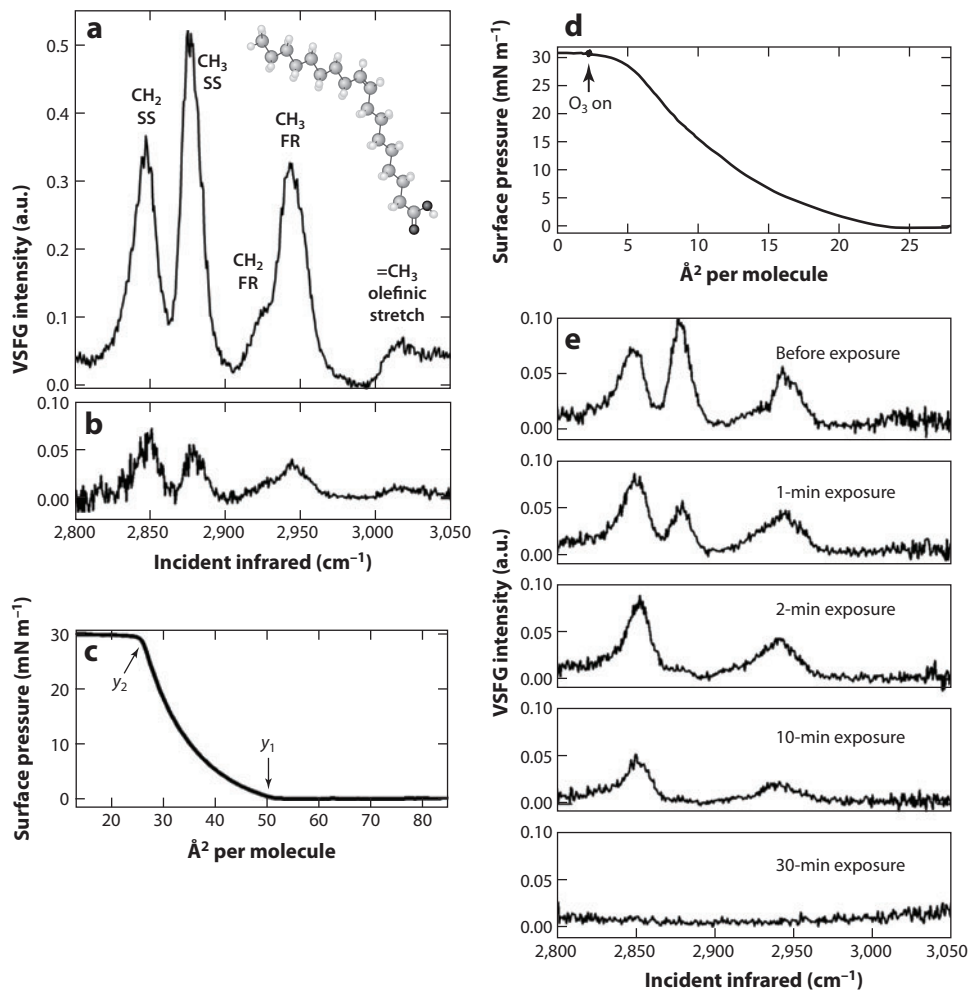
for the PA monolayer results from the slight deprotonation of the PA head group at pH 6, which produces a surface charge. Surface charges have been shown to enhance water structure through charge-dipole interactions for a variety of systems at both the vapor/water and solid/water interfaces (31, 32, 35, 70, 87, 114, 130). These enhancements have also been widely attributed to  $\chi^{(3)}$  effects, as first reported by Zhao et al. (114) for water near charged lipid monolayers. The inset of **Figure 6a** shows a physical representation of this. For PA monolayers spread on  $\text{CaCl}_2$ -containing subphases (**Figure 6c**), the degree of enhancement observed for the peaks in the O-H region is less than that observed for PA on a neat water subphase, even for weakly concentrated (0.1-M)  $\text{CaCl}_2$  solutions. This is further evidence of the ion complexation that occurs between the  $\text{COO}^-$  head group and  $\text{Ca}^{2+}$  ion in solution. Ion complexation effectively screens the surface charge induced by the deprotonated PA head group, as predicted by Gouy-Chapman theory (114, 130–133). The screening effect is seen to increase with  $\text{Ca}^{2+}$  concentration, as expected, until the water spectrum for a PA monolayer spread on a 1.8-M  $\text{CaCl}_2$  subphase (**Figure 6c**) resembles the water spectrum for a 1.8-M  $\text{CaCl}_2$  solution without a PA monolayer at the vapor/solution interface (**Figure 6d**), except for the disappearance of the peak at  $\sim 3,700\text{ cm}^{-1}$  (86). The disappearance of the  $3,600\text{ cm}^{-1}$  peak in the spectra corresponding to the PA monolayer on the 1.8-M  $\text{CaCl}_2$  subphase is the final confirmation that the head group of the associated PA molecules is completely deprotonated, as this mode has been attributed to the O-H mode of the protonated PA head group (32, 86).

Recently PS-SFG has been applied to the study of the water structure near charged and neutral Langmuir monolayers on neat water subphases (31–33, 35, 87, 123). These studies generally demonstrate the reorientation of water molecules causes the O-H transition dipole to point either toward or away from the charged monolayer, depending on the sign of the charge, although our recent work on zwitterions is an exception to this trend, as these act as negatively charged head groups (32). Nihonyanagi et al. (87) recently extended this method to study the water structure near charged lipid monolayers spread on  $\text{NaCl}$ -containing subphases and observed similar decreases in O-H peak intensity on the addition of salts, as is observed for a PA SAM spread on ion-containing subphases. Although Nihonyanagi et al. invoked electrolyte screening of a Gouy-Chapman electric double layer by the added salts and not an explicit ion-complexation argument, their results highlight the suitability of the PS-SFG technique to investigate complex lipid/ion/water systems.

These results demonstrate the intricate interactions that can occur for fat-coated marine aerosols between the surface organic film and the aqueous subphase. Interactions such as deprotonation have consequences for the growth and uptake of fat-coated aqueous aerosols as deprotonation can induce packing in the monolayer by reducing the charge repulsion of the head groups via ion complexation, with ions contained within the subphase such as  $\text{Mg}^{2+}$  and  $\text{Ca}^{2+}$  (74, 86, 88). Tightly packed monolayers are generally more resistant to water uptake because of the increased hydrophobicity of the outer layer, thus inhibiting further growth of the aerosol, although this idea has been challenged (105, 134, 135).

## 4.2. Oxidation of Oleic Acid Monolayers

Several researchers have recently applied VSFG to the study of the oxidation, generally by gaseous ozone ( $\text{O}_3$ ), of tropospherically relevant organic films. Stokes and colleagues (136–138) used VSFG to monitor the reaction and reaction products of glass slides functionalized with tropospherically relevant organic moieties with ozone at the vapor/solid interface. Their results indicated reaction probabilities for a variety of organic moieties with ozone that are consistent with the molecular dynamics simulations of Vieceli et al. (139). At the vapor/liquid interface, more relevant to aqueous phase aerosols, Voss et al. (73) monitored the reaction for an OA monolayer on an aqueous



**Figure 7**

Conventional ssp polarized VSG spectra and surface pressure–area isotherms ( $\pi$ - $A$ ) of oleic acid (OA) monolayers on neat water subphases. (a) VSG spectrum of an OA monolayer on neat water at a surface pressure of  $\sim 12$  mN m $^{-1}$ . Abbreviations: FR, Fermi resonance; SS, symmetric stretch. (Inset) Molecular model of the OA molecule. (b) VSG spectrum of an OA monolayer at a surface pressure of  $\sim 1$  mN m $^{-1}$ . (c) Langmuir compression isotherm with markers indicating the point of film collapse ( $y_2$ ) and the point at which the VSG spectrum of panel b ( $y_1$ ) was acquired. (d) Langmuir isotherm during oxidation of oleic acid monolayer by ozone. (e) VSG spectra corresponding to the oxidation of OA monolayers at the air/water interface with ozone after set exposure times. Figure reprinted with permission from Reference 73. Copyright 2007 American Geophysical Union.

subphase with ozone via VSG. The oxidation of OA has been widely used as a proxy for the heterogeneous reaction of fat-coated aerosols with atmospheric oxidants because of its prevalence in atmospheric particulate matter (112, 115, 122, 140).

Here we focus the discussion on the oxidation for OA monolayers at the vapor/water interface by ozone in conjunction with the above discussion on the organization and reaction of atmospherically relevant fatty acid lipid monolayers. **Figure 7a** presents the VSG spectrum of an OA



monolayer spread on a neat water subphase compressed to a surface pressure of  $\sim 12$  mN m<sup>-1</sup>, which corresponds to an average area per OA molecule in the film of 32 Å per molecule, slightly before the monolayer collapses at 28 Å per molecule indicated by point  $y_2$  in **Figure 7c**. There are five spectral features present in **Figure 7a** that are attributed to the CH<sub>2</sub> symmetric stretch (2,846 cm<sup>-1</sup>), the CH<sub>3</sub> symmetric stretch (2,876 cm<sup>-1</sup>), the CH<sub>2</sub> Fermi resonance (2,923 cm<sup>-1</sup>), the CH<sub>3</sub> Fermi resonance (2,941 cm<sup>-1</sup>), and the olefinic = CH stretch (3,014 cm<sup>-1</sup>) (73). **Figure 7b** illustrates the VSFG spectrum for an OA monolayer slightly compressed to  $\sim 1$  mN m<sup>-1</sup>, indicated by point  $y_1$  in **Figure 7c**. Unlike for the PA monolayers discussed above, the CH<sub>2</sub> vibrational modes of the OA monolayer possess strong VSFG intensity owing to a lack of inversion for the CH<sub>2</sub> moieties, which results from the unsaturated nature of the OA molecule.

Spectra corresponding to the reaction of the OA film with ozone are shown in **Figure 7e**, and the corresponding Langmuir isotherm during the oxidation is shown in **Figure 7d**. Here an immediate drop in the CH<sub>3</sub> symmetric stretch peak is apparent after 1 min of exposure to ozone, and by 30 min of exposure, all VSFG signatures of the OA monolayer are gone. Through separate VSFG studies on the known reaction products of OA with ozone (140)—nonanal [CH<sub>3</sub>(CH<sub>2</sub>)<sub>7</sub>CHO], nonanoic acid [CH<sub>3</sub>(CH<sub>2</sub>)<sub>7</sub>CO<sub>2</sub>H], azelaic acid [HO<sub>2</sub>C(CH<sub>2</sub>)<sub>7</sub>CO<sub>2</sub>H], and 9-oxononanoic acid [OCH(CH<sub>2</sub>)<sub>7</sub>CO<sub>2</sub>H]—it was determined that the spectra shown in **Figure 7e** do not originate from reaction products at the vapor/water surface, which is consistent with the relatively high solubilities for these reaction products in water (141). To explain the observed spectral changes for the oxidation of OA, Voss et al. (73) compared the oxidation spectra (**Figure 7e**) with the VSFG spectra observed for a slightly compressed ( $\sim 1$  mN m<sup>-1</sup>), and thus disordered, OA monolayer at the vapor/water interface (**Figure 7b**). The similarities between the spectra shown in **Figure 7b** and the spectra corresponding to the short ozone exposure times (1-min exposure) indicated that oxidation disordered the OA film. These results suggest that the oxidation of an organic film on an aqueous aerosol induces disorder within this film via the dissolution of the oxidation products into the aqueous subphase. This creates space for the unoxidized OA molecules to disperse owing to steric effects, analogous with observations for the slightly compressed OA monolayer (**Figure 7b**) (73). This disorder may lead to the fragmentation of the aerosol or to heightened evaporation of the aqueous subphase; both these scenarios would lead to the formation of smaller fat-coated aerosols.

## 5. FUTURE PERSPECTIVES FOR VSFG OF ENVIRONMENTAL INTERFACES

Many recent advances in instrumentation and methodology have made the future of applying VSFG to the study of environmental interfacial processes a bright one. Perhaps the most exciting is the recent development of the PS-SFG technique (31), which enables the direct measurement of the transition dipole moment and provides insight into the organization of molecules. However, the application of this technique has been limited by the small number of research groups adopting it (31–33, 35), likely because of the complexity of the instrumentation and spectral processing. As an alternative, the utilization of the maximum entropy method has shown promise (18, 19, 34).

Recent advances have also been made for time-resolved VSFG measurements (142–144) as well as nonlinear optical spectroscopic measurements (145–148) from nonplanar interfaces such as colloidal suspensions. The ability to resolve dynamics at interfaces is a crucial step toward furthering the understanding of how molecules behave within the constrained interfacial region. Colloidal suspensions are ubiquitous in the environment and extremely important for the aqueous geochemistry of pollutant cycling (4).

## SUMMARY POINTS

1. VSFG is a powerful tool for the investigation of environmental interfaces. Water molecules at the vapor/solution interface are shown to reorganize in the presence of ions. The size, charge, geometry, and polarizability of ions play a role in this reorganization, making PS-SFG direct measurements of the interfacial water structure important.
2. Solutes and lipids display very different behavior at the vapor/water interface than ions because of their increased surface preference. The orientation of these species at the surface can be determined through VSFG.
3. Reactions involving ions, solutes, and lipids in interfacial regimes are observable with VSFG. This is demonstrated by observing a series of various reactions for each class of systems: ion pairing between nitrate and the counterion at the vapor/solution interface, the reaction and uptake of methanol by sulfuric-acid solutions, and ion complexation and oxidation for various lipids spread on aqueous subphases.

## DISCLOSURE STATEMENT

The authors are not aware of any affiliations, memberships, funding, or financial holdings that might be perceived as affecting the objectivity of this review.

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7. Demonstrates for the first time the experimental power of nonlinear optical phenomena.

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11. Observes for the first time the vapor/neat water interface with VSFG.

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25. Summarizes early VSFG contributions to understanding molecular behavior at the vapor/solution interface for ionic solutions and solutes.

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31. Demonstrates a generally applicable method for the experimental determination of the imaginary component of the second-order molecular susceptibility.

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46. Observes the increase in interfacial depth upon salt addition at the vapor/solution interface via VSFG.

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47. Observes the presence of hydronium ions in the vapor/solution interface for HCl, HBr, and HI solutions.

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## Errata

An online log of corrections to Annual Review of Physical Chemistry chapters (if any, 1997 to the present) may be found at <http://physchem.AnnualReviews.org/errata.shtml>