
Neutralization of Methyl Cation via Chemical Reactions in Low-Energy Ion-Surface Collisions with Fluorocarbon and Hydrocarbon Self-Assembled Monolayer Films

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Low-energy ion-surface collisions of methyl cation at hydrocarbon and fluorocarbon self-assembled monolayer (SAM) surfaces produce extensive neutralization of CH_3^+ . These experimental observations are reported together with the results obtained for ion-surface collisions with the molecular ions of benzene, styrene, 3-fluorobenzonitrile, 1,3,5-triazine, and ammonia on the same surfaces. For comparison, low-energy gas-phase collisions of CD_3^+ and 3-fluorobenzonitrile molecular ions with neutral *n*-butane reagent gas were conducted in a triple quadrupole (QQQ) instrument. Relevant MP2 6-31G*/MP2 6-31G* *ab initio* and thermochemical calculations provide further insight in the neutralization mechanisms of methyl cation. The data suggest that neutralization of methyl cation with hydrocarbon and fluorocarbon SAMs occurs by concerted chemical reactions, i.e., that neutralization of the projectile occurs not only by a direct electron transfer from the surface but also by formation of a *neutral* molecule. The calculations indicate that the following products can be formed by *exothermic* processes and without appreciable activation energy: CH_4 (formal hydride ion addition) and C_2H_6 (formal methyl anion addition) from a hydrocarbon surface and CH_3F (formal fluoride addition) from a fluorocarbon surface. The results also demonstrate that, in some cases, simple thermochemical calculations cannot be used to predict the energy profiles because relatively large activation energies can be associated with exothermic reactions, as was found for the formation of CH_3CF_3 (formal addition of trifluoromethyl anion). (J Am Soc Mass Spectrom 2002, 13, 1151–1161) © 2002 American Society for Mass Spectrometry

Low-energy collisions of polyatomic ions with surfaces and the reaction mechanisms that characterize them represent an ongoing research challenge. Detailed summaries and descriptions of ion-surface collision processes have been provided in reviews [1–3] and in recent papers by, e.g., Cooks and coworkers [4] and Hayward et al. [5]. Briefly, these processes can be classified into at least five general categories: Elastic scattering, neutralization, surface-induced dissociation, ion-surface atom/group transfer reactions, and chemical sputtering [4]. Reactions in each category can occur via multiple mechanisms, and often it is difficult to distinguish the various pathways [4, 5]. The remarkably large number of proposed mechanisms provides evidence of the vast multitude of possibilities [1–9].

Ion neutralization at surfaces is quite different from the other four processes because (1) neutralization leads to neutral species that are not detected *directly* by mass spectrometry, and (2) it results in a loss of ion signal that obviously affects sensitivity. There is a general consensus regarding the overall neutralization process: Neutralization of a positively charged projectile occurs via electron transfer from the surface. Although the electron-transfer mechanism is difficult to describe (for attempts, see, e.g., Reference [7] and references therein), it is supported by several experimental observations, including surface-induced dissociation (SID) efficiency measurements [1, 2, 7–9] and thermochemical and quantum chemical calculations [9]. It should be noted that ion-surface collisions of atomic projectiles with bare metal surfaces [10] must be distinguished from collisions between small organic (polyatomic) ions and a metal surface covered with an ordered (long chain) self-assembled organic monolayer (SAM). In the case of organic ions colliding with an organic surface, which is discussed in this paper, the neutralization is less exten-

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sive than that for atom/bare metal collisions. The amount of neutralization at self-assembled monolayer surfaces is roughly dependent on the energy difference between the ionization energy of the organic surface species and the recombination energy of the projectile. The amount of neutralization at fluoroalkanethiolate and alkanethiolate SAM surfaces varies, for example, in the order $\text{Ar}^+ > \text{Cl}^+ > \text{CH}_3^+ > [\text{benzene}]^+ > \text{Na}^+ (\text{K}^+)$ [9]. Increased ion signal (less neutralization) observed on a fluoroalkanethiolate surface relative to an alkanethiolate surface has been related to the higher ionization energy of fluorinated versus hydrogenated alkyl chains [2, 8, 9]. Other related studies on CF_3 -terminated Langmuir-Blodgett films [11] and SAMs showed that significantly less neutralization of small organic projectile ions occurs on fluorocarbon surfaces than on hydrocarbon surfaces. These studies, which utilized SAMs such as $\text{CF}_3(\text{CH}_2)_{15}\text{S-Au}$ and $\text{CF}_3\text{CF}_2(\text{CH}_2)_{14}\text{S-Au}$ [12] and projectiles such as benzene and pyrazine, show that a single CF_3 terminal group (or a CF_3CF_2 group) at the uppermost layer of these surfaces mimics a more highly fluorinated surface (e.g., $\text{CF}_3(\text{CF}_2)_9\text{CH}_2\text{CH}_2\text{S-Au}$). Note that SID experiments using isotopically labeled Langmuir-Blodgett films also led to the conclusion that predominantly the uppermost atoms or groups are involved in low-energy ion-surface reactions [13].

In the present work, the role of the terminal groups of self-assembled monolayers in ion neutralization is explored by *ab initio* and thermochemical calculations together with experimental observations associated with the neutralization of methyl cation at several distinct SAM surfaces. The studies show that neutralization of methyl cation can occur via *concerted chemical reactions*, i.e., not only by *direct* electron transfer from the surface. (Of course, an electron must be transferred from the metal substrate to neutralize the ionized surface chain; otherwise, the surface would be charged, which has not been observed for SAMs used in our SID studies [14].) Specifically, for CH_3^+ , the following products can be formed in *exothermic* processes and without appreciable activation energy: CH_4 (hydride ion addition), C_2H_6 (methyl anion addition), and CH_3F (fluoride addition). The studies also demonstrate that, in some cases, simple thermochemical calculations cannot be reliably used to predict the energy profiles, given that relatively large activation energy can be associated with exothermic reactions, as was found for the formation of CH_3CF_3 (trifluoromethyl anion addition).

Experimental and Computational

In experimental studies, we used perdeutero methyl cation, CD_3^+ , which was generated from CD_3I or acetone- d_6 (Aldrich, Milwaukee, WI) by electron impact (EI) ionization (70 eV). No difference was observed when CD_3^+ was generated from either CD_3I or acetone- d_6 . Benzene, styrene, 3-fluorobenzonitrile, ammonia, and 1,3,5-triazine molecular ions were generated similarly from the corresponding neutral molecules (Al-

drich). The alkanethiols investigated include alkyl-terminated [$\text{CH}_3(\text{CH}_2)_{15}\text{SH}$, denoted by C16F0], CF_3 -terminated [$\text{CF}_3(\text{CH}_2)_{15}\text{SH}$, denoted by C16F1], CF_3CF_2 -terminated [$\text{CF}_3\text{CF}_2(\text{CH}_2)_{14}\text{SH}$, denoted by C16F2], and $\text{C}_{10}\text{F}_{21}$ -terminated [$\text{CF}_3(\text{CF}_2)_9(\text{CH}_2)_6\text{SH}$, denoted by C16F10]. The shortened nomenclature C16Fx indicates total number of carbons (16) and the number of carbons from the alkyl terminus (x) that are perfluorinated. The normal alkanethiol, $\text{CH}_3(\text{CH}_2)_{15}\text{SH}$ (92%), was purchased from Aldrich and used without further purification. The synthesis of the fluorocarbon-terminated alkanethiols has been reported elsewhere [12, 15].

Vapor-deposited gold surfaces were purchased from Evaporated Metal Films (Ithaca, NY). The gold surfaces were formed on a 17 mm \times 13 mm \times 0.5 mm silica base that was covered with a 5 nm adhesion under-layer of titanium and then 100 nm of vapor-deposited gold. The gold surfaces were UV-cleaned for 15 min (Boeckel UV cleaner, Model 135500) before being placed in appropriate thiol solutions. Ethanolic stock solutions of the various alkanethiols (~ 1 mM) were prepared in glassware cleaned with piranha solution ($\text{H}_2\text{SO}_4/\text{H}_2\text{O}_2$, in a 3:1 volume ratio; *Caution: Piranha solution should be handled with extreme care!*). A 6–8 mL aliquot of the ethanolic stock solution of a particular thiol was transferred to a new disposable test tube after the tube was UV-cleaned for 5 min. Monolayer formation occurred by allowing the freshly UV-cleaned gold surface to react in the ethanol thiol solution for 72 h. The films were then rinsed six times by sonication in ethanol, dried under argon, and inserted into the fast entry lock of the mass spectrometer (placed under vacuum) within 5–10 min.

A detailed description of the basic SID tandem mass spectrometer with recent modifications has been published elsewhere [16, 17]. Briefly, two Extrel quadrupole mass analyzers are arranged in a 90° configuration with an adjustable surface holder bisecting the mass analyzers at a 45° angle. Collision energy is controlled by the potential difference between the ion source and the surface, where mass-selected singly charged ions from the first quadrupole collide into the organic thin film. The products of the collision are then focused into and analyzed by the second quadrupole of the tandem system. Ion focusing is accomplished by a set of three metal lenses placed after the first quadrupole and in front of the second quadrupole (focus ions to and from the surface). The surface holder secures four surfaces, and each surface can be individually placed in the ion beam path of the mass spectrometer. This setup allows all four surfaces to be analyzed under the same experimental conditions. During these experiments, the base pressure of the instrument was 5×10^{-7} torr and sample pressure rarely exceeded 9×10^{-7} torr (measured in the analysis chamber).

For neutralization studies, currents were measured with a picoammeter at the target surface during the collision. Elevated current at the target surface during collision indicates an enhanced electron transfer that

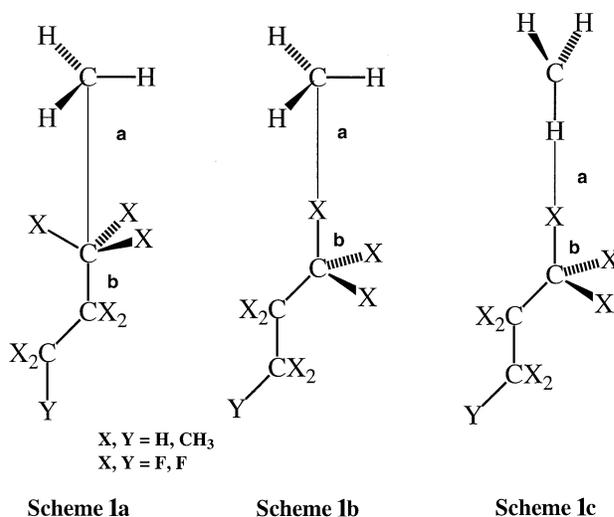
neutralizes the incoming projectile ions (correspondingly, enhanced surface neutralization leads to low ion current measured at the electron multiplier of the mass spectrometer). The measurement of the extent of neutralization was the comparison of the current generated from the collision of the projectile ion beam with an UV-cleaned bare gold surface. The bare gold surface was UV-cleaned for 15 min and placed in absolute ethanol until inserted into the mass spectrometer. The neutralization of the projectile ion beam with a bare metal surface is assumed to be ~99% efficient and independent of the projectile ion. Therefore, determining the amount of ions that are neutralized and survive the collision is expressed as the neutralized ion beam (NIB) and the surviving ion beam (SIB): % NIB = (current for SAMs)(99%)/(current for bare metal); % SIB = 100 - % NIB.

Although the absolute neutralization efficiency at a bare Au surface is not known exactly, 99% neutralization is a reasonable assumption considering early results of SID experiments on bare metal surfaces and disordered short chain alkanethiols [1, 8]. More importantly, slight deviations from the assumed 99% value do not affect the relative SIB values for all projectile ions reported in this paper.

In our preliminary studies, external electric fields caused drift in current measurements that limited our ability to collect reproducible data. The precision of the current measurements was improved with the addition of a Faraday cage. A wire mesh (Faraday) cage was placed around the dual quadrupole SID instrument on all four sides and on the top. The cage was grounded to the vacuum chamber during the current measurements.

Ion/molecule reactions were performed in a Finnigan TSQ 700 triple quadrupole mass spectrometer (Finnigan MAT, San Jose, CA). Product ion mass spectra were recorded by mass selecting an ion of interest using the first quadrupole and allowing the ion to pass through the second quadrupole (rf-only), which contained a reagent gas. The third quadrupole was scanned over the mass range of interest at a rate of one second/scan. The ion source and vacuum manifold were held at 150 °C and 70 °C, respectively. The collision cell offset potential was set to produce a low-energy laboratory collision of 4 eV. The reagent gas used for reaction was *n*-butane (+99%, Aldrich), which was introduced directly into the collision quadrupole by linking the *n*-butane lecture bottle regulator to the collision cell gas port on the TSQ. The pressure of the reagent gas was controlled with the TSQ leak valve and maintained at ~1.4 mTorr.

The closed shell methyl cation is sufficiently small for quantum chemical calculations. Unfortunately, even one surface chain of the SAM is too large for high-level *ab initio* calculations on our university processors. However, the similarity of SID and gas-phase collision experiments [8, 18] and reactions occurring between [benzene]⁺ cation and ¹³CH₃⁻, CD₃⁻, and CF₃-terminated Langmuir-Blodgett [11, 13] and self-assembled monolayer films [12] indicate that primarily the upper-



most atoms or group of a surface chain are involved in low-energy ion-surface reactions. These observations justify the use of smaller compounds to model the surface chain. We used *n*-C₄H₁₀ to model an alkanethiolate surface and C₃F₈ to model a fluorinated alkanethiolate surface (Scheme 1). [For comparison purposes, specifically for the addition of CF₃ (trifluoromethyl anion), we also used CH₃CH₂CH₂CF₃ to model the CF₃-terminated CF₃(CH₂)₁₅S-Au surface.] The experimental observation of charged reaction products indicates that the projectile ion slows down and stays at the vicinity of the surface (2–3 Å) for *at least* one vibrational period (ca. 10⁻¹³ s). This feature justifies the use of the Born-Oppenheimer approximation, i.e., these reactions can be modeled by the time-independent Schrödinger equation. Strictly speaking, our computational model below assumes that projectile ions have low kinetic energy in close vicinity of the surface during the neutralization. Of course, a percentage of higher energy ions can exist and can exit the potential energy surface (PES) at higher energy pathways.

Although in real ion-surface collisions a large number of random orientations are possible, the three specific orientations illustrated in Schemes 1a, b, and c were considered in the present work. As will be shown, these orientations are reasonable models for describing certain characteristic features of ion-surface reactions and projectile neutralization. All calculations were performed at the MP2 6-31G* level by using the program package Gaussian 94 [19]. The potential energy surfaces were calculated by changing the -CX₃...CH₃ distance and the -CX₂...CX₃ distance (distances a and b in Scheme 1a, respectively). All other geometry parameters were optimized assuming the retention of C_s symmetry during the model reaction. Note again that the C_s symmetry restriction may introduce some artifacts in the potential energy surfaces. However, we believe that such an arbitrarily chosen symmetry restriction is a better approximation for a "real" ion-surface reaction because the rotation of the "reacting" chain (or upper-

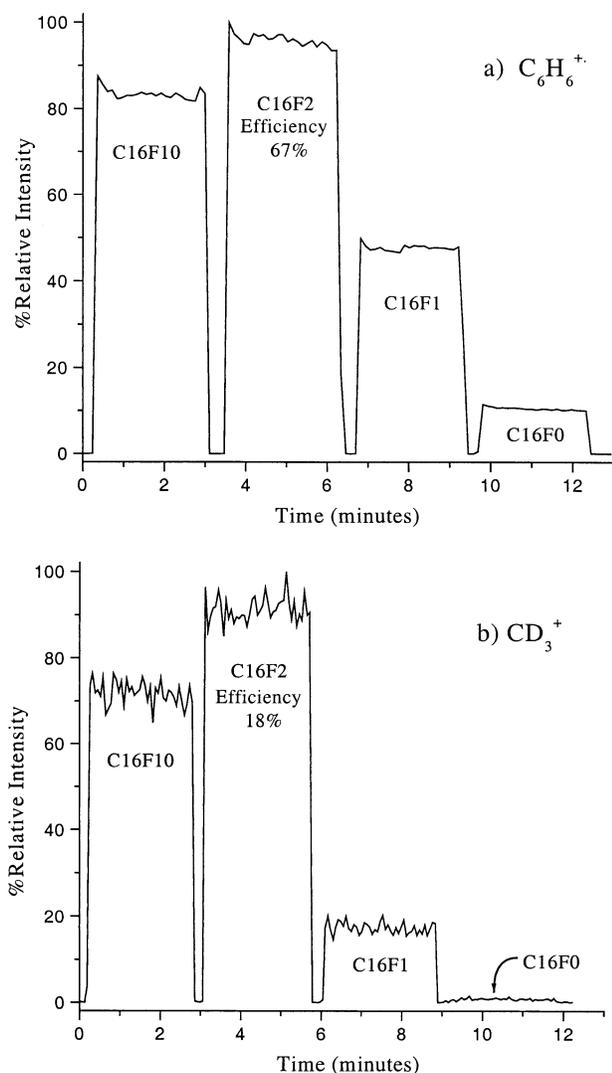


Figure 1. The total scattered-ion current (TIC) was measured at the detector of the SID instrument for (a) benzene radical cation ($[benzene]^+$) and (b) perdeutero methyl cation (CD_3^+) collided at 30 eV kinetic energy into $C_{10}F_{21}$ -terminated (C16F10), CF_3CF_2 -terminated (C16F2), CF_3 -terminated (C16F1), and alkyl-terminated (C16F0) SAM films. The time axis represents the actual acquisition time of the ion-surface collision experiments, i.e., ~3 min of acquisition time for each surface.

most group) is hindered by neighboring chains (or groups) in a well-ordered SAM film.

Results and Discussion

Experimental Results

Figures 1a and b show the total scattered-ion current (TIC) recorded by the detector of the SID tandem mass spectrometer when $[benzene]^+$ and CD_3^+ cations, respectively, were collided at 30 eV SID energy with four distinct SAM surfaces. The corresponding 30 eV product ion spectra recorded for the CD_3^+ cation at two surfaces (C16F2 and C16F0) are shown in Figure 2.

The comparison of Figure 1a and Figure 1b indicates

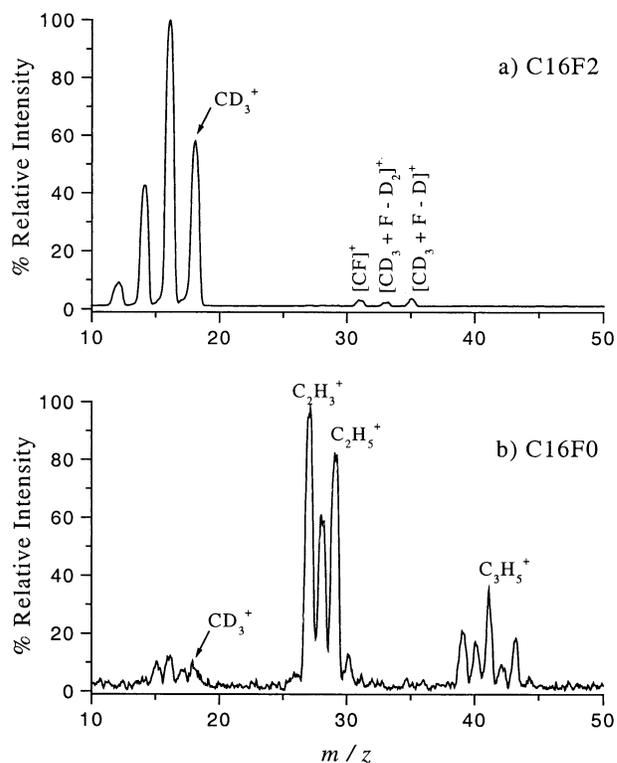


Figure 2. Product spectra formed by collision between perdeutero methyl cation (CD_3^+) at 30 eV kinetic energy with (a) CF_3CF_2 -terminated (C16F2) and (b) alkyl-terminated (C16F0) SAM films.

both the similarities and differences between the total scattered-ion current of $[benzene]^+$ and CD_3^+ cations with the same hydrocarbon and fluorocarbon surfaces. The highest ion recovery (i.e., the least amount of neutralization) was observed for both cations at the CF_3CF_2 -terminated surface, C16F2 (Figure 1). A closer look at the TIC results in Figure 1 reveals important differences between the neutralization of $[benzene]^+$ and CD_3^+ . Although a general trend exists for the two cations with the hydrocarbon and fluorocarbon surfaces, the measured SIB percentage values are characteristically different. (For the definition of % SIB, see the Experimental and Computational section). For $[benzene]^+$, ~67% SIB was measured at the C16F2 surface, which is in remarkably good agreement with the SID efficiency (I_{ev}/I_0) results published earlier for a C12F10 surface [8]. The corresponding % SIB value obtained at the same surface for the CD_3^+ cation is only ~18%. These data were collected three independent times with the present instrumental setup and are highly reproducible. Consistently, the signal-to-noise (S/N) ratio is much better in the benzene product-ion spectra (spectra not shown for $[benzene]^+$ but also manifested in the lower variability of signal intensity on the top of the TIC curves in Figure 1a versus 1b).

Interestingly, 30 eV collisions with the "more fluorinated" surface (C16F10) resulted in slightly lower TIC current for both projectiles than obtained for the C16F2 surface. This small difference may be related to slight

differences in the conformation and packing density of these surfaces. In particular, recent analyses by atomic force microscopy (AFM) have found that the lattice spacing of the CF_3CF_2 -terminated SAM is ~ 5.0 Å (i.e., indistinguishable from that of normal alkyl-terminated SAMs on gold) and the lattice spacing of the $\text{C}_{10}\text{F}_{21}$ -terminated SAM is ~ 5.8 Å [20–23].

The C16F1 surface is also interesting because it exhibits different TIC behavior toward $[\text{benzene}]^+$ versus CD_3^+ cations. If the % SIB value for each projectile with the C16F1 surface is normalized to the C16F2 surface % SIB value, the normalized % SIB ratio (C16F2: C16F1) for benzene is 1.0:0.62 while the same ratio for the methyl cation is 1.0:0.27. Previous ion-surface collisions studies with pyrazine molecular ion with the C16F1 surface have shown that hydrogen atoms (from the CH_2 group under the CF_3 group) are exposed and available for reaction [12]. Therefore, the origin of the TIC difference (increase of neutralization of methyl cation) with the C16F1 surface is at least in part from the smaller size of CD_3^+ versus $[\text{benzene}]^+$. The steric difference permits more facile reaction of the methyl cation with the second-layer (highly reactive) CH_2 groups leading to the formation of neutral reaction products (as previously mentioned and explained in detail below). Although more neutralization (lower TIC) is observed at the C16F1 surface (with only a CF_3 -terminal group) than for the C16F10 and C16F2 surfaces for both projectiles, it is also clear from Figure 1 that the inclusion of only one CF_3 terminal group characteristically improves ion survival compared to normal alkyl-terminated surfaces (C16F0).

Another notable difference between the TICs of $[\text{benzene}]^+$ and CD_3^+ is that the TIC obtained at the hydrocarbon surface (C16F0) is only slightly above baseline for CD_3^+ , but clearly above baseline for $[\text{benzene}]^+$. Although it was difficult to measure the % SIB for the CD_3^+ cation at the C16F0 alkanethiolate surface, it is unquestionably less than 1.0%. Accordingly, the CD_3^+ cation is virtually missing from the product ion spectrum, and mostly sputtered ions (alkyl chain fragments from the surface) are detected at 30 eV SID collision energy (see Figure 2b). On the other hand, the S/N ratio is better in the product ion spectrum of CD_3^+ obtained with the C16F2 partially fluorinated surface (Figure 2a). In this spectrum, CD_3^+ and its fragments are the dominant ions. The presence of ions at m/z 31, 33, and 35 indicates fluoride addition and subsequent fragmentation of the reaction product (see labels in Figure 2a). The F-addition reaction products were verified with ion-surface collisions utilizing the unlabeled methyl cation (CH_3^+). However, the presence of the ion at m/z 31 (CF^+) could be a combination of fragmented reaction product as well as sputtering from the fluorocarbon SAM surface.

It is important to note again that the % SIB values for CD_3^+ have been found to be characteristically and consistently lower than those obtained for $[\text{benzene}]^+$ on the same surface. The % SIB values are: $18 \pm 2\%$ for

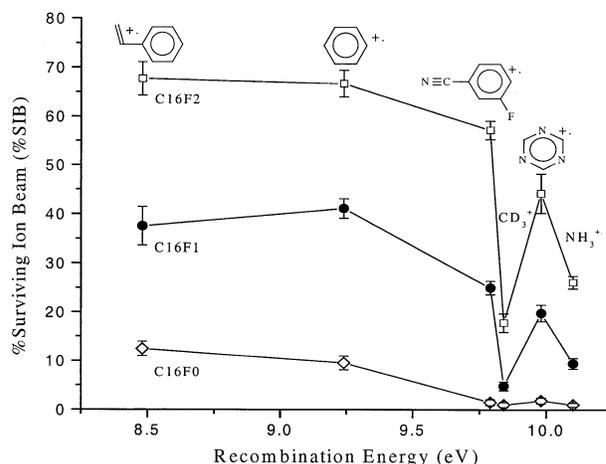


Figure 3. Percentage surviving ion beam (% SIB) as a function of the recombination energy of selected projectile ions. These results were generated by collisions between the projectile ions at 20 eV SID energy with (a) CF_3CF_2 -terminated (C16F2: open squares), (b) CF_3 -terminated (C16F1: closed circles), and (c) alkyl-terminated (C16F0: open diamonds) SAM surfaces.

CD_3^+ and $67 \pm 3\%$ for C_6H_6^+ on the fluorocarbon (C16F2) surface; $<1.0\%$ for CD_3^+ and $9.5 \pm 1.4\%$ for C_6H_6^+ on the hydrocarbon (C16F0) surface. The large difference in SID efficiency for CD_3^+ versus $[\text{benzene}]^+$ is surprising given the relatively small difference between the ionization energies [24] of CH_3 (9.84 eV) and benzene (9.25 eV). Based on a simple electron-transfer mechanism, i.e., relying simply on the difference between the ionization energy of the SAM surface and the recombination energy of the projectile, the neutralization of CH_3^+ at a fluorinated surfaces is predicted to be strongly endothermic while the same process at alkanethiolate surfaces is weakly endothermic. These simple calculations, which will be discussed below in greater detail, indicate that processes other than simple electron transfer might play a role in the neutralization of CH_3^+ at fluorocarbon and hydrocarbon SAM surfaces.

To further explore the role of the SAM ionization energy and the recombination energy of the projectile, we performed low energy ion-surface collision experiments for four additional projectile ions having recombination energies (RE) that lie in the range of ca. 8.5–10.1 eV [24]. (Note that, strictly speaking, the recombination energies are not known directly for all ions of interest; therefore, the ionization energies of the corresponding neutrals (molecules or radicals) are assumed to be the recombination energies in this paper. By convention, we use positive RE values, even though ion recombination is related to energy release.) These projectiles include the molecular ions of styrene (RE = 8.45 eV), 3-fluorobenzonitrile (RE = 9.78 eV), 1,3,5-triazine (RE = 9.95 eV), and NH_3 (RE = 10.10 eV). The % SIB values as a function of recombination energies for three different SAM surfaces (C16F2, C16F1, and C16F0) are shown in Figure 3. The general trend is a decreasing % SIB with increasing recombination energy. The CD_3^+ cation con-

sistently exhibits much lower % SIB values than would be expected given its recombination energy. For example, although the recombination energy of the 3-fluorobenzonitrile molecular ion is almost the same as that of CD_3^+ (9.78 eV versus 9.84 eV), there is a significant difference in the % SIB values obtained for these cations. We selected the NH_3^+ cation to mimic the geometry and size of the small CH_3^+ cation, although one can argue that the difference in electron parity (i.e., the open electron structure NH_3^+ and the closed shell CH_3^+) might give rise to different neutralization processes. In a recent related study [25], the surviving ion beam (% SIB) results generated with oxygen-containing projectiles (1,4-dioxane: RE = 9.19 eV and acetone: RE = 9.70 eV) on the same surfaces have shown ion-neutralization behavior that deviates from the trend in Figure 3 (e.g., C16F2: % SIB ~44% for 1,4-dioxane and ~43% for acetone). Taken as a whole, these data suggest that ion neutralization is dependent on the atomic composition and/or structure of the projectile ions. The % SIB results for oxygen-containing projectiles are still substantially higher when compared to the results for methyl cation. In any case, the % SIB values obtained for the projectiles presented in Figure 3 support the assumption that processes other than a "simple" electron transfer might also take place during CH_3^+ neutralization at SAM surfaces.

As was mentioned in the Experimental section, we also conducted certain low energy ion-molecule reactions in a Finnigan triple quadrupole instrument. Figure 4a shows the resulting spectrum from ion-molecule reactions between (neutral) *n*-butane reagent gas and CD_3^+ . The base peak in this spectrum is the butyl cation at m/z 57. The low intensity peak at m/z 18 corresponds to the CD_3^+ cation. No fragmentation of CD_3^+ is observed, which is not surprising given that the collision cell offset potential was set to produce a 4 eV laboratory collision in these experiments. At higher collision energies (up to 9 eV), more m/z 18 survives the collision (the m/z 18/57 ratio increases) suggesting that low collision energies favor a reactive long-lived collision complex that allows the formal hydride transfer [26]. Some low intensity butane fragment ions (m/z 43, 41, and 29) also appear in the spectrum. These fragments mimic the sputtered ions in the SID experiments described above. In the 70 eV EI spectrum of $\text{C}_4\text{H}_{10}^+$, the intensity of the fragment ion at m/z 57 is much lower than that of the molecular ion (spectrum not shown). We therefore conclude that the ion at m/z 57 originates predominantly from a specific hybrid transfer ion-molecule reaction that is associated with the formation of neutral methane [26]: $\text{CD}_3^+ + \text{C}_4\text{H}_{10} \rightarrow \text{CD}_3\text{H} + \text{C}_4\text{H}_9^+$.

This experiment demonstrates that this reaction is preferred in the gas-phase and suggests that an analogous reaction can also occur at alkanethiolate self-assembled monolayer surfaces (see above).

For comparison, we also conducted an ion-molecule reaction between neutral *n*-butane reagent gas and the molecular ion of 3-fluorobenzonitrile (RE = 9.78 eV) in

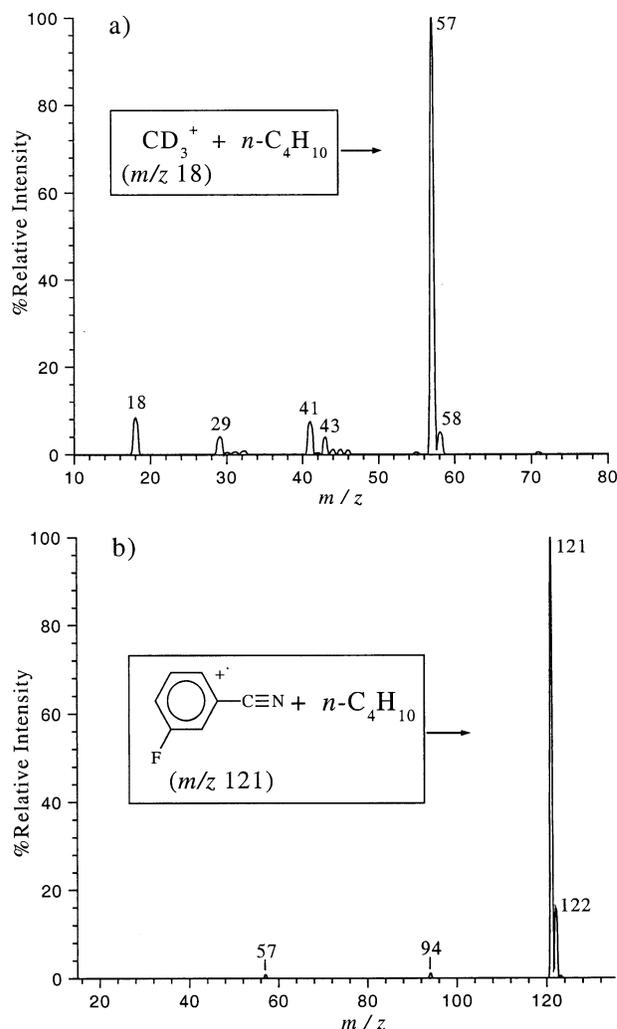


Figure 4. Spectra generated from ion-molecule reactions between *n*-butane reagent gas and (a) perdeutero methyl cation (m/z 18) and (b) 3-fluorobenzonitrile ion (m/z 121) in a Finnigan triple quadrupole instrument. In both cases, the collision cell offset potential was set to produce a low-energy laboratory collision of 4 eV.

the triple quadrupole instrument. The resulting spectrum is shown in Figure 4b. In strong contrast to the spectrum of Figure 4a, the intensity of the peak at m/z 57 is extremely low (<2%), and the base peak is the 3-fluorobenzonitrile molecular ion (m/z 121). These results are consistent with the SID results discussed above: Although the recombination energies of CD_3^+ and 3-fluorobenzonitrile are nearly indistinguishable (9.84 eV versus 9.79 eV) [24], much less neutralization is observed for the 3-fluorobenzonitrile ion than for the CD_3^+ ion.

Computational Results

To provide theoretical support for some of the possible alternative neutralization processes of CH_3^+ at these surfaces, we performed a series of MP2 6-31G*//MP2 6-31G* ab initio calculations. Three different orienta-

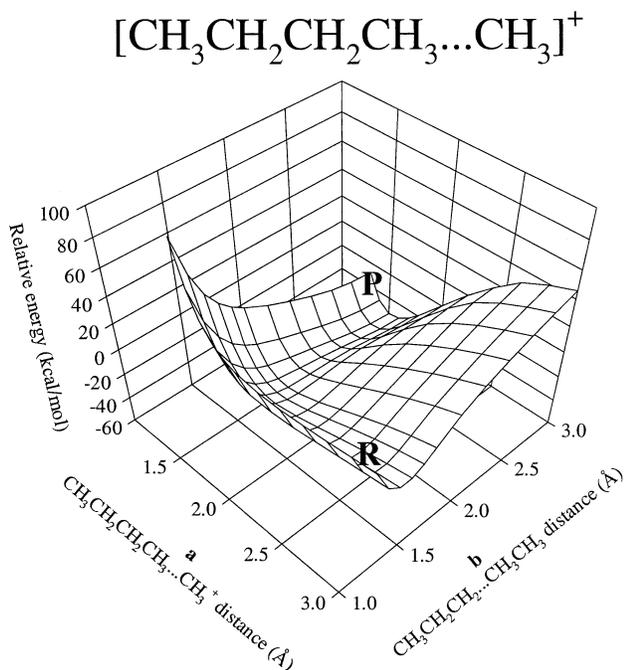


Figure 5. MP2 6-31G**/MP2 6-31G* potential energy surface was generated for a model collision between CH_3^+ and $n\text{-C}_4\text{H}_{10}$. Parameters a and b in Scheme 1a were varied over the range of 1.3–3.0 Å. R indicates the reactants (CH_3^+ and C_4H_{10}), and P indicates the products ($\text{CH}_3\text{CH}_3 + \text{C}_3\text{H}_7^+$).

tions were modeled as shown in Scheme 1a, b, and c. In all cases, the CH_3^+ approaches the surface (represented by distance a). Preliminary calculations [9] showed that as the CH_3^+ approaches the surface with the “chain” orientations shown in Scheme 1a and b, the $\text{CH}_2\text{-CH}_3$ carbon bond of the alkanethiolate model surface (Scheme 1a) and the C–H bond of the uppermost methyl group of the surface (Scheme 1b) are significantly elongated. Therefore, in the present work, we used these bond lengths (represented by b in Schemes 1a and b) as a second variable parameter to obtain two-dimensional potential energy surfaces. With a similar approach, we used the $\text{CH}_3\text{CH}_2\text{CH}_2\text{CH}_2\dots\text{H}$ distance (distance b in Scheme 1c) to model a head-on collision.

Figure 5 shows the MP2 6-31G**/MP2 6-31G* potential energy surface obtained for the model reaction between methyl cation and $\text{CH}_3\text{CH}_2\text{CH}_2\text{CH}_3$ in the “reactive” configuration shown in Scheme 1a, (X = H, Y = CH_3). This reaction models the formation of neutral CH_3CH_3 by the methyl cation (abstraction of $-\text{CH}_3$ from the surface chain). The lower left axis of Figure 5 indicates the distance between the CH_3^+ projectile and the surface (distance a in Scheme 1a), while the right axis (distance b in Scheme 1a) indicates the $\text{CH}_2\text{-CH}_3$ carbon-carbon bond length of the cleaving alkyl chain. The reference energy (0 kcal/mol) is the sum of the MP2 6-31G**/MP2 6-31G* total energy of $n\text{-C}_4\text{H}_{10}$ and CH_3^+ . The reference energy is the energy of the system at infinite surface-projectile distance and is virtually un-

changed until CH_3^+ approaches the surface closer than ca. 4 Å. (In other words, ca. 4 Å is the distance at which the methyl cation starts to “see” or “feel” the surface.) An important feature of the potential energy surface in Figure 5 is that the system is rolling downhill, i.e., the methyl anion addition reaction proceeds with no activation barrier to produce neutral ethane and a charged surface (C_3H_7^+ in our model, see product side, P, in Figure 5). Although the number of C–C and C–H bonds is the same in both the reactants and products, the exothermicity of this reaction is due to the increased “delocalization” of the positive charge in C_3H_7^+ relative to CH_3^+ . Unfortunately, the neutral product, ethane, goes undetected by the mass spectrometer, but the measured low SID efficiency value for CH_3^+ (i.e., the loss of ~99% of CH_3^+ ion current) is consistent with the theoretical prediction. Strictly speaking, the potential energy surface in Figure 5 is valid if there is no “spontaneous” or “concurrent” rearrangement(s) of the ionized surface chain. Although we have no evidence to exclude such a rearrangement, intuitively one can assume that it would require a longer time, especially for highly packed SAM chains so that the formed neutral (CH_3CH_3) can already be far away from the surface.

It should be noted that in the SID spectrum of CH_3^+ obtained on the hexadecanethiolate (C_{16}F_0) surface (Figure 2b), chemical sputtering peaks are seen almost exclusively. Sputtered ions originating from the surface are also good indicators of projectile neutralization because a prerequisite for their formation is a transfer of the charge from the projectile to the surface (together with a significant internal energy transfer to the surface to cleave the C–C bonds of the surface chains). The formation of the sputter peaks could originate from direct electron transfer and/or hydride ion transfer with the methyl ion (similar to the fragment peaks observed from the C_4H_{10} during the ion-molecule reactions experiment). The presence of sputtered peaks in Figure 1b could support a direct electron-transfer ($\text{CH}_3^+ + \text{C}_4\text{H}_{10} \rightarrow \text{CH}_3 + \text{C}_4\text{H}_{10}^+$) argument and may contradict the formation of neutral ethane (or methane, see below) predicted by the ab initio calculations.

Note, however, that in contrast to neutralization of CH_3^+ via the formation of ethane, direct electron transfer is a slightly endothermic process (Figure 6). Based on experimental heats of formation for the reactants and products [24], the reaction enthalpies for the formation of ethane ($\text{CH}_3^+ + \text{C}_4\text{H}_{10} \rightarrow \text{CH}_3\text{CH}_3 + \text{C}_3\text{H}_7^+$) and the direct electron transfer (from $n\text{-C}_4\text{H}_{10}$ to CH_3^+ producing CH_3 and $n\text{-C}_4\text{H}_{10}^+$) are -40.2 and 16.7 kcal/mol, respectively (Figure 6). Note that the exothermicity of the reaction leading to CH_3CH_3 and C_3H_7^+ may be underestimated if the rearrangement of C_3H_7^+ (SAM chain) is possible. We must emphasize that we do not claim that all CH_3^+ cation is neutralized via the formation of ethane (or methane, see below). With reference to the reaction enthalpies discussed above (Figure 6) and the lack of an ab initio reaction barrier, we conclude that, particularly at low collision energies, the neutral-

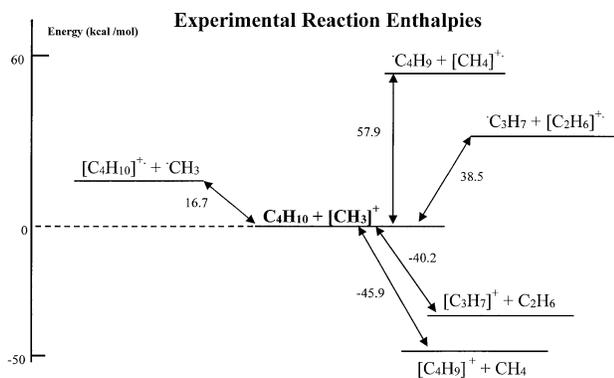


Figure 6. Reaction enthalpies of some possible reactions between CH_3^+ and $n\text{-C}_4\text{H}_{10}$. All values are based on experimental heats of formation taken from Reference [24].

ization of methyl cation via the formation of neutral ethane (or methane) can compete with direct electron transfer from $n\text{-C}_4\text{H}_{10}$ to CH_3^+ . (At higher collision energies, however, the energy of the real system might lie above the calculated potential energy surface shown in Figure 5; consequently, this surface will not be useful in predicting processes that occur at higher energies.)

Although not shown here, a MP2 6-31G*//MP2 6-31G* potential energy surface (similar to the formation of ethane, Figure 5) was obtained for the hydride

ion addition reaction that leads to the formation of CH_4 and C_4H_9^+ (Scheme 1b, $X = \text{H}$, $Y = \text{CH}_3$). Again, this reaction is calculated to have no activation barrier, which suggests that the formation of CH_4 is an alternative reaction pathway for neutralizing methyl cation. The hydride ion addition reaction is also predicted to be exothermic based on experimental heats of formation (-45.9 kcal/mol, Figure 6). Its existence is supported by the ion-molecule results (intense formation of butyl ion, $m/z = 57$, in Figure 4a) obtained using CD_3^+ collisions with neutral n -butane gas in which the major detected product is C_4H_9^+ (presumably $\text{CD}_3^+ + \text{C}_4\text{H}_{10} \rightarrow \text{C}_4\text{H}_9^+ + \text{CD}_3\text{H}$). Intuitively, hydride ion addition to form methane should be sterically favored over the abstraction of methyl anion to form ethane.

The "head on" collision (i.e., $\text{H} \dots \text{H}$, Scheme 1c) is shown as (1) a MP2 6-31G*//MP2 6-31G* potential energy surface as well as (2) a contour graph for better visualization in Figure 7. The potential energy surface is highly repulsive (notice the absolute values of the relative energies at short a and b distances in Scheme 1c), which indicates energy deposition to *both* the projectile and the surface (kinetic energy/potential energy, T-V, conversion). Two pathways from reactants (R) to products (P1 or P2) are illustrated in Figure 7. Path 1 (P1) shows a collision in which the surface C–H bond

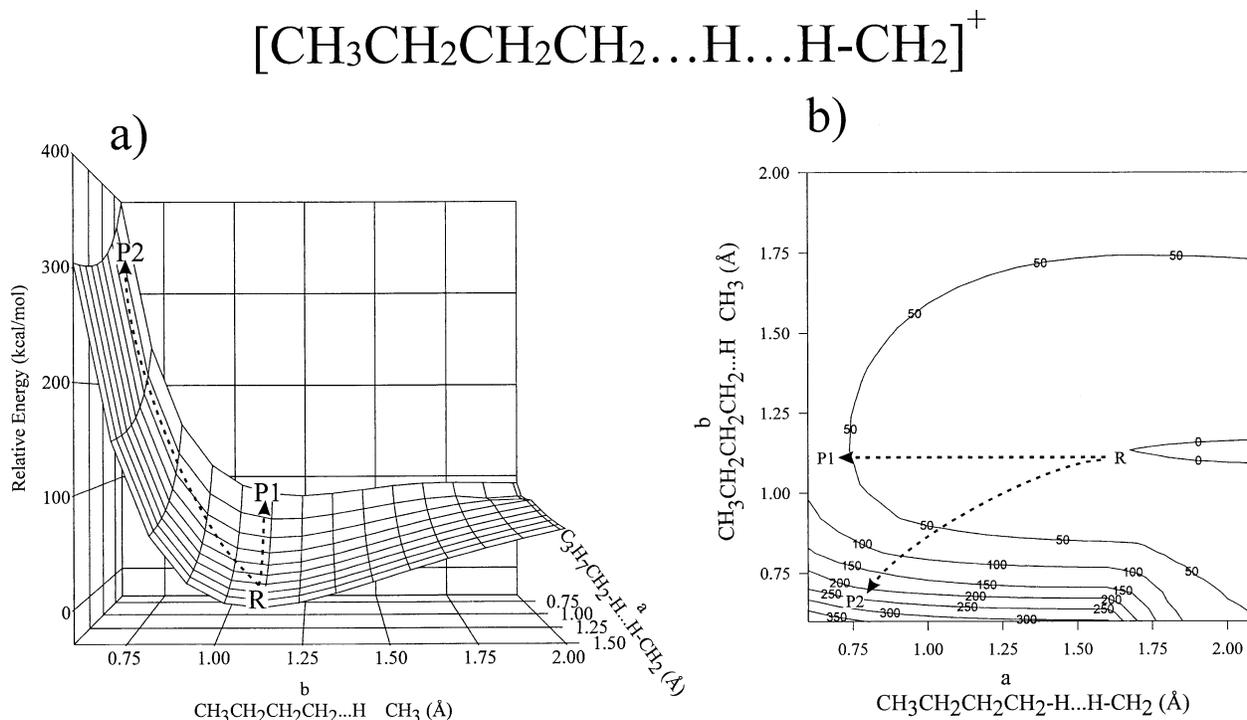


Figure 7. (a) MP2 6-31G*//MP2 6-31G* potential energy surface and (b) contour graph was generated for a "head on" ($\text{H} \dots \text{H}$) collision between CH_3^+ and $n\text{-C}_4\text{H}_{10}$. Parameters a and b in Scheme 1c were varied over the range of 0.75–3.0 Å. Two pathways going from reactants (R) to products (P1 or P2) are illustrated in Figure 7 as reaction pathways going from reactants (R) to products (P1 or P2). Path 1 (P1) shows a collision in which the surface C–H bond (distance b) is held constant at ca. 1.1 Å. Path 2 (P2) shows when the surface C–H bonds are distorted during the collision.

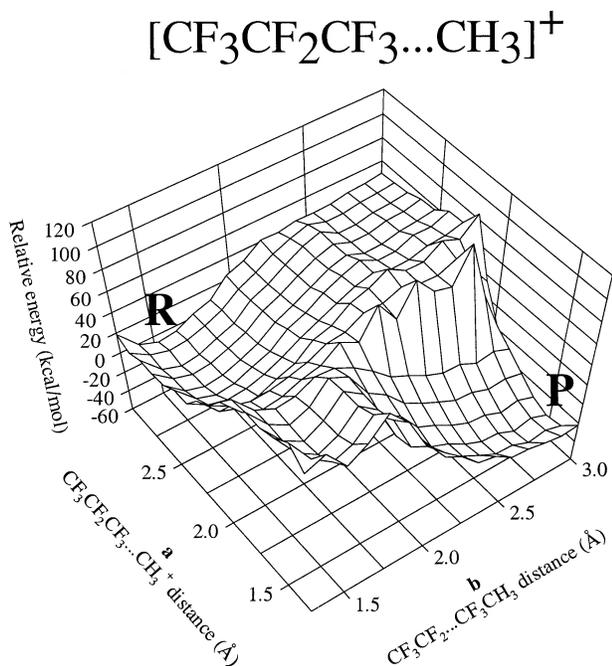


Figure 8. MP2 6-31G*//MP2 6-31G* potential energy surface for a model collision between CH_3^+ and C_3F_8 . Parameters a and b in Scheme 1a were varied over the range of 1.3–3.0 Å. R indicates the reactants (CH_3^+ and C_3F_8), and P indicates the products (CH_3CF_3 + C_2F_5^+).

(distance b) is held constant at ca. 1.1 Å. With this rigid surface, the relative energy increases as the methyl cation approaches the surface (decrease in distance a). In reality, the surface bonds are also distorted during the collision, and the energy of the system is highly repulsive (Path 2) and can even exceed 10 eV (>240 kcal/mol). The calculations also show that the “head on” collision between H...H atoms (Scheme 1c) fails to give neutralization (~60% of the charge still remains on the methyl cation even at a H...H distance of 0.6 Å.) These results, together with those discussed above, suggest that the outcome of low-energy ion-surface reactions (e.g., neutralization versus energy deposition) depends on the relative orientation of the projectile and the outermost group of the surface.

Figure 8 shows the MP2 6-31G*//MP2 6-31G* potential energy surface calculated for the reaction, $\text{CH}_3^+ + \text{C}_3\text{F}_8 \rightarrow \text{CH}_3\text{CF}_3 + \text{CF}_2\text{CF}_3^+$, which is the “fluorinated” analog of the ethane formation described above. The orientation is shown in Scheme 1a ($X = \text{F}$, $Y = \text{F}$), and both parameters a and b were varied over the range of 1.3–3.0 Å. In Figure 8, the left axis represents the approach of CH_3^+ to the fluorocarbon chain (parameter a in Scheme 1a), and the right axis represents the cleavage of the $\text{CF}_3\text{CF}_2\text{...CF}_3$ bond (parameter b in Scheme 1a). In contrast to Figure 5, the potential energy surface in Figure 8 indicates that the addition of CF_3 (trifluoromethyl anion), i.e., the neutralization of CH_3^+ by the formation of CH_3CF_3 , possesses a substantial activation barrier (ca. 60 kcal/mol). It is important to

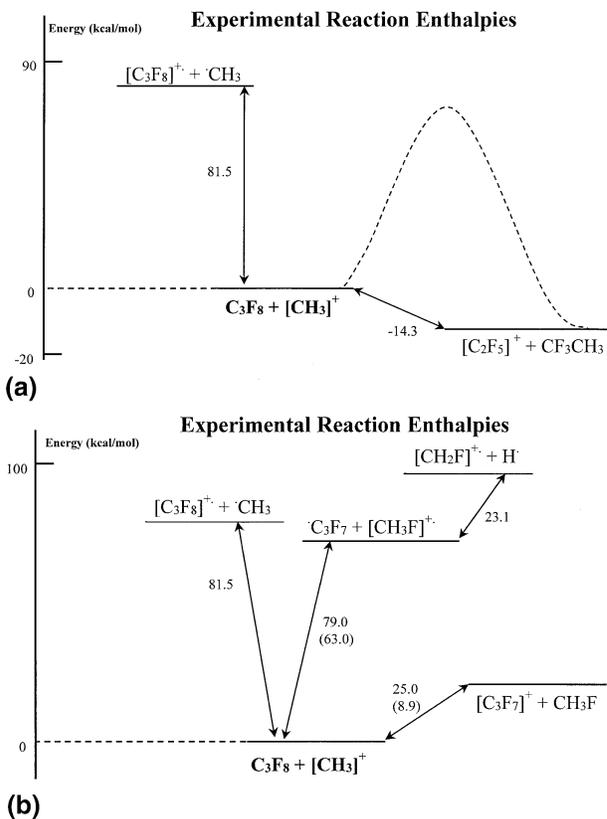


Figure 9. (a) Comparison between reaction enthalpies of a charge-exchange process and a CF_3^- (trifluoromethyl anion) addition process for the collision between CH_3^+ and C_3F_8 . All values are based on experimental heats of formation taken from Reference [24]. (b) Reaction enthalpies calculated for some possible reactions between C_3F_8 and CH_3^+ (0 kcal/mol reference energy). All values are based on experimental heats of formation taken from Reference [24]. Two values are shown for the reactions $\text{C}_3\text{F}_7 + [\text{CH}_3\text{F}]^+$ and $[\text{C}_3\text{F}_7]^+ + \text{CH}_3\text{F}$, which are based on two different estimations of the heats of formation for C_3F_7^+ (–96.8 kcal/mol [24c] and –80.8 kcal/mol [24c]) where the former value corresponds to the value in the parentheses.

note that the reaction $\text{CH}_3^+ + \text{C}_3\text{F}_8 \rightarrow \text{CH}_3\text{CF}_3 + \text{CF}_2\text{CF}_3^+$ is predicted to be *exothermic* by simple thermochemical calculations using experimental heats of formation ($\Delta H_r = -14.3$ kcal/mol, Figure 9a) [24], which highlights an important limitation of using “simple” thermochemical calculations.

A similar energy surface with a similar barrier has been calculated for the reaction $\text{CH}_3^+ + \text{CF}_3\text{CH}_2\text{CH}_2\text{CH}_3 \rightarrow \text{CH}_3\text{CF}_3 + \text{CH}_2\text{CH}_2\text{CH}_3^+$; the results (potential surface not shown, but it is similar to energy surface shown in Figure 8) demonstrate the importance of the terminal CF_3 group. The origin of the barrier for addition of $-\text{CF}_3$ to CH_3^+ (from C_3F_8 or $\text{CH}_3\text{CH}_2\text{CH}_2\text{CF}_3$) can be associated with the lone pairs of the F atoms of the outermost CF_3 group of a fluorinated chain, i.e., to molecular orbitals that are formed by the lone pairs of F atoms. We can visualize these orbitals as “umbrella” or “shield” orbitals that prevent the “penetration” of the empty orbital of the methyl cation in line with the $\text{CF}_2\text{–CF}_3$ bond. Because of the

lack of such an interaction, neither the direct neutralization (“flow” of electrons from the surface to the projectile) nor the addition of CF_3 (trifluoromethyl anion) is a preferred process. Such an interaction between appropriately aligned carbon orbitals is unhindered when hydrogen atoms are present (as in regular hydrocarbons), which explains why calculations show that *both* hydride ion and methyl anion addition occur. These additions allow extensive neutralization of methyl cation with alkanethiolate surfaces [2, 8, 9]. Of course, a strictly linear arrangement is unnecessary for the reactions with CH_3^+ ; a slight “off-axis” approach (i.e., when the α, β angle deviates from 180 degrees) can also lead to similar products. The importance of orbital interactions has been recognized, for example, for ion reactions at clean and oxygen-modified Ni (111) surfaces [7].

No energy barrier was predicted at the MP2 6-31G*//MP2 6-31G* level for the addition of fluoride, i.e., for the reaction $\text{CH}_3^+ + \text{CF}_3\text{CF}_2\text{CF}_3 \rightarrow \text{CH}_3\text{F} + \text{CF}_2\text{CF}_2\text{CF}_3^+$ (potential energy surface not shown, but it is similar to that shown in Figure 5). The favored formation of CH_3F can rationalize why methyl cation survival is much lower (ca. 12%) than benzene survival (ca. 68%) on the fluorinated surface (see Figure 1). As shown in Figure 9b, experimental heats of formation predict the reaction $\text{CH}_3^+ + \text{CF}_3\text{CF}_2\text{CF}_3 \rightarrow \text{CH}_3\text{F} + \text{CF}_2\text{CF}_2\text{CF}_3^+$ to be slightly endothermic although there is an ambiguity in the heats of formation of C_3F_7 and C_3F_7^+ [24]. However, the analogous reaction, $\text{CH}_3^+ + \text{CF}_3\text{CF}_3 \rightarrow \text{CH}_3\text{F} + \text{CF}_2\text{CF}_3^+$, is thermoneutral ($\Delta H_r = 0.7$ kcal/mol). For the molecular ion of benzene, a similar interaction is sterically more hindered, i.e., the “reactive configuration” between the benzene molecular ion and the outermost CF_3 group of the fluorinated surface is statistically less favorable than other orientations that lead to T–V conversion. Nevertheless, the reactive orientation is still accessible. The addition of fluorine to benzene molecular ion has been observed in SID spectra on fluorinated surfaces at relatively low abundance [8], but the neutralization of $[\text{benzene}]^+$ is not associated with this reaction.

Conclusions

The experimental and computational results presented in this paper indicate that neutralization of CH_3^+ cation can occur by two types of processes: (1) Direct electron transfer from the SAM surface and (2) concerted chemical reactions without activation barriers. The decreasing neutralization of methyl cation on fluorocarbon surfaces relative to hydrocarbon surfaces can be attributed to the presence of lone pairs on the F atoms of the terminal CF_3 groups. The lone pairs inhibit the direct electron-transfer process (from the SAM to CH_3^+) and the direct formation of CH_3CF_3 . The computational results also indicate that the relative orientation of the projectile influences the “outcome” of the modeled ion-surface reactions. Be-

cause most of the SAM (and/or Langmuir-Blodgett) surfaces used in SID studies are well ordered, as characterized by various techniques such as AFM, STM, and wettability, their “static” behavior is independent of the projectile ion. The “dynamic” behavior, however, is the result of the interaction of the projectile ion and the outermost layer of the surface chains. Several ion-surface group orientations can lead to different processes, such as (1) “T–V” orientations, in which energy deposition is preferred (see, head-on collision model, Scheme 1c), and (2) “neutralization” or “reactive” orientations in which interactions between orbitals of the ion and the surface are allowed. These orbital interactions can then lead to either simple electron transfer or chemical reactions, the products of which can be neutrals (see, examples for the formation of CH_3CH_3 , CH_4 , and CH_3F) or charged species (see, e.g., fluorine, F, addition in Figure 2). From the point of view of the projectile, we can describe a “reactivity” (or “neutralization”) volume, which is a part of space in the vicinity of the ion within which a reaction with a reactive partner is allowed. Although this volume is “dynamic” rather than “static” in character, one can intuitively assume that the smaller the projectile, the larger this volume (with reference to the projectile size). For methyl cation, for example, this “reactive” volume can be defined as a “cone” above and below the plane of the cation. For a much larger projectile, e.g., for a protonated peptide, only small and local “volumes” can be reactive but are sterically hindered by the atoms of groups of the projectile. It is not surprising, therefore, that no ion-surface reactions have been reported for protonated peptides. Note that while “odd electron” molecular ions of common organic compounds are reported to be reactive, the “odd electron” character is not a prerequisite for reactivity (see our present example, the methyl cation).

Even if they are carried out on small model systems, theoretical calculations can be helpful in describing and predicting the relative contributions of the different processes that dominate low-energy ion-surface collisions (for example, molecular dynamics modeling of energy deposition of various projectile ions at diamond and SAM surfaces [27, 28]) and might also contribute to the development of more sophisticated models that can be used to describe low-energy ion-surface reactions.

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